

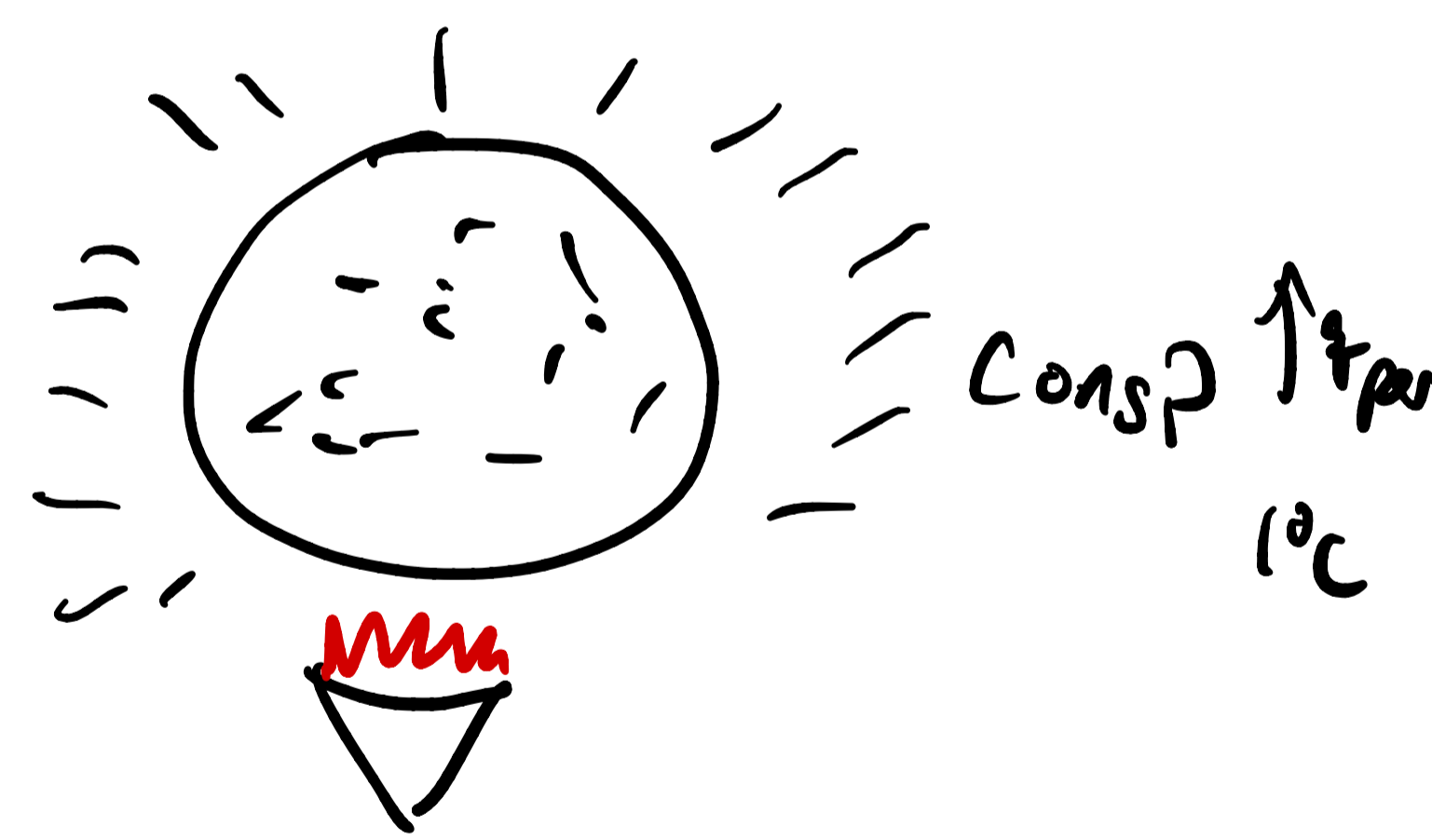
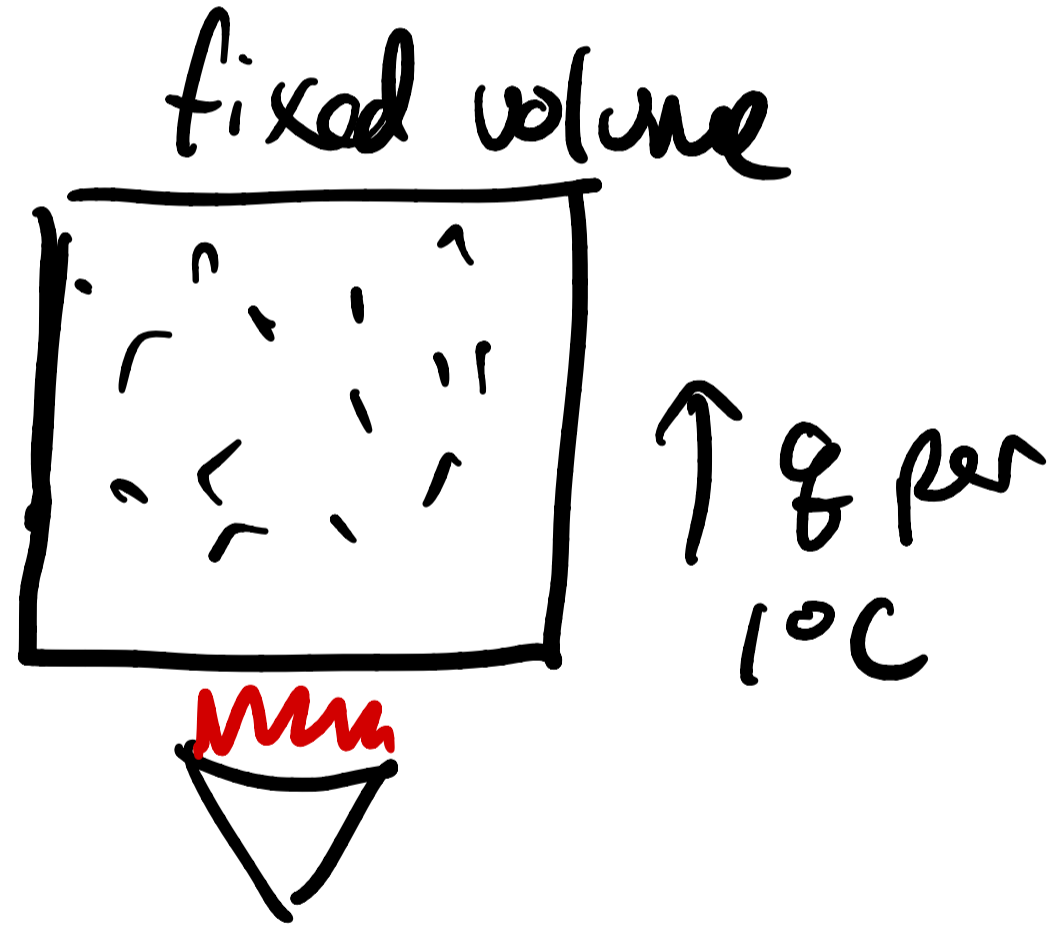
Lecture 6

heat capacity (extensive quantity)

$$dq = C dT$$

$$C_v = \left(\frac{\partial q}{\partial T} \right)_v$$

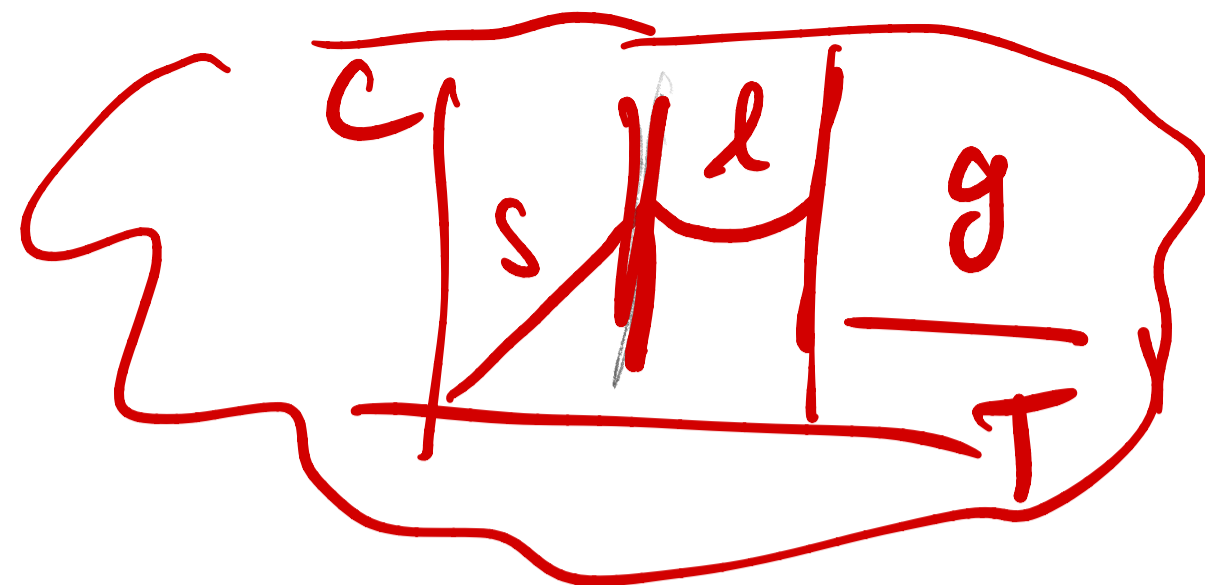
$$C_p = \left(\frac{\partial q}{\partial T} \right)_p$$



$$C_p > C_v$$

How much heat flows:

$$dq = C(T) dT$$



$$q = \int_{T_1}^{T_2} C(T) dT \approx C \int_{T_1}^{T_2} dT = C(T_2 - T_1)$$

if T_1 & T_2 aren't that diff

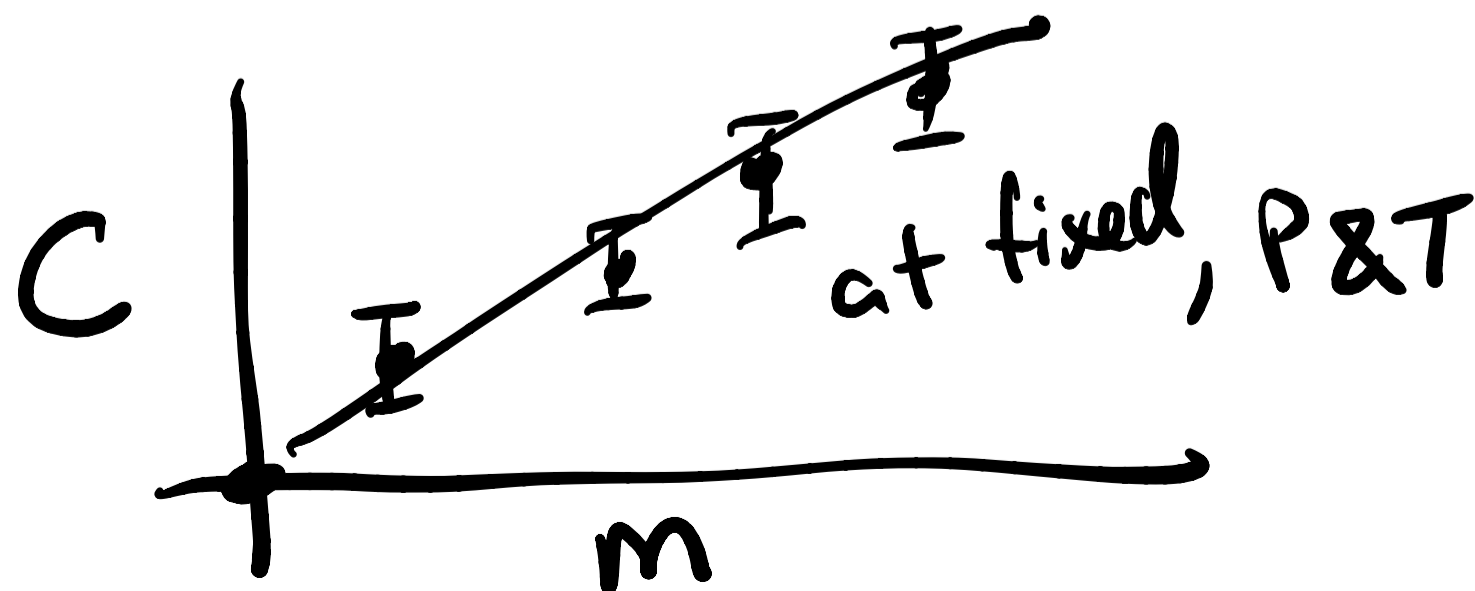
& phase transition, $C(T) \approx C$

→ " $q = n C \Delta T$ "

\hat{L} molar

" $q = m C \Delta T$ "

\hat{L} per mass



How much does E change

$$dE = dq + dW \quad (\text{1st law})$$

$$= C dT - P dV$$

$$\Delta E = \int_a^b (C dT - P dV)$$

a) const volume, $dV = 0$ $C = C_V$

$$\Delta E = Q \Delta T$$

b) const P , $C = C_P$

$$\Delta E = C_P \Delta T - P \Delta V$$

$$C_v = \left(\frac{\partial \mathcal{E}}{\partial T} \right)_v = C_v \frac{\partial \mathcal{E}}{\partial T} = C_v$$

$$d\mathcal{E} = C_v dT$$

$$dq = d\mathcal{E}$$

(a) C_v is how much energy changes
for some change in temp

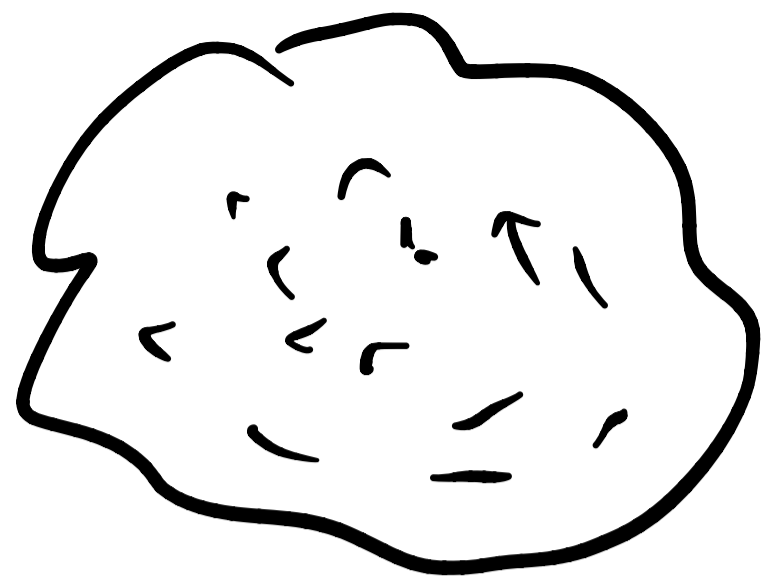
b) $d\mathcal{E} = dq - PdV \Rightarrow dq = d\mathcal{E} + PdV$

$$C_p = \left(\frac{d(\mathcal{E} + PdV)}{dT} \right)_p$$

$$d\mathcal{E} + PdV$$

Enthalpy

$$H \equiv E + PV$$



$$d(H) = d(E) + d(PV)$$

$$= dE + VdP + PdV$$

Consider $dP = 0$

@ const P

$$dH = dE + PdV = dq$$

$$C_P = \left(\frac{\partial H}{\partial T} \right)_P$$

Ideal gas: (monatomic ideal gas)

2 assumptions: - take up no space
- don't interact

$$(1) PV = nRT$$

$$(2) E = \frac{3}{2} nRT$$

$$E_{\text{total}} = KE^{\propto T} + \cancel{PE}^{\circ}$$

$$C_V^{\text{ig}} = \left(\frac{\partial E}{\partial T} \right)_V = \frac{3}{2} nR$$

$$C_V = \frac{C_V}{n} = \frac{3}{2} R$$

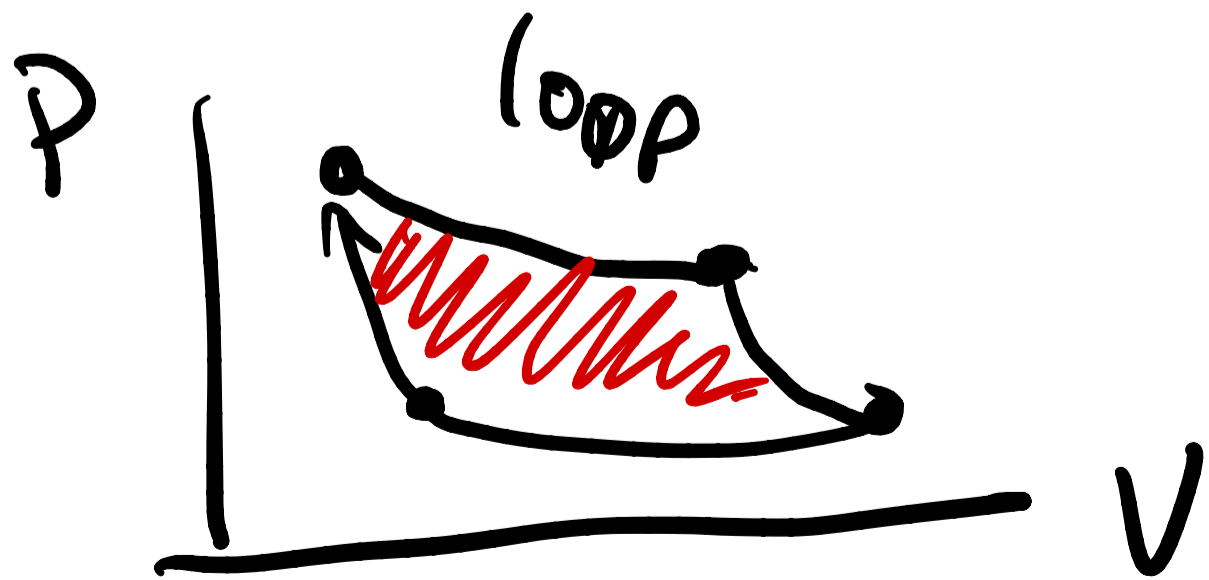
$$C_P^{\text{ig}} = \left(\frac{\partial (E + PV)}{\partial T} \right)_P = \frac{\partial}{\partial T} \left(\frac{3}{2} nRT + nRT \right) = \frac{5}{2} nR$$

$C_p > C_v$, ideal gas

$$C_p - C_v = nR$$

→ true for any # atoms

Preview : how to do work
through a cyclic process



$$W = \int_{\text{cycle}} -P dV$$

Change state:

- Const P , $dP=0$

- Const V , $dV=0$

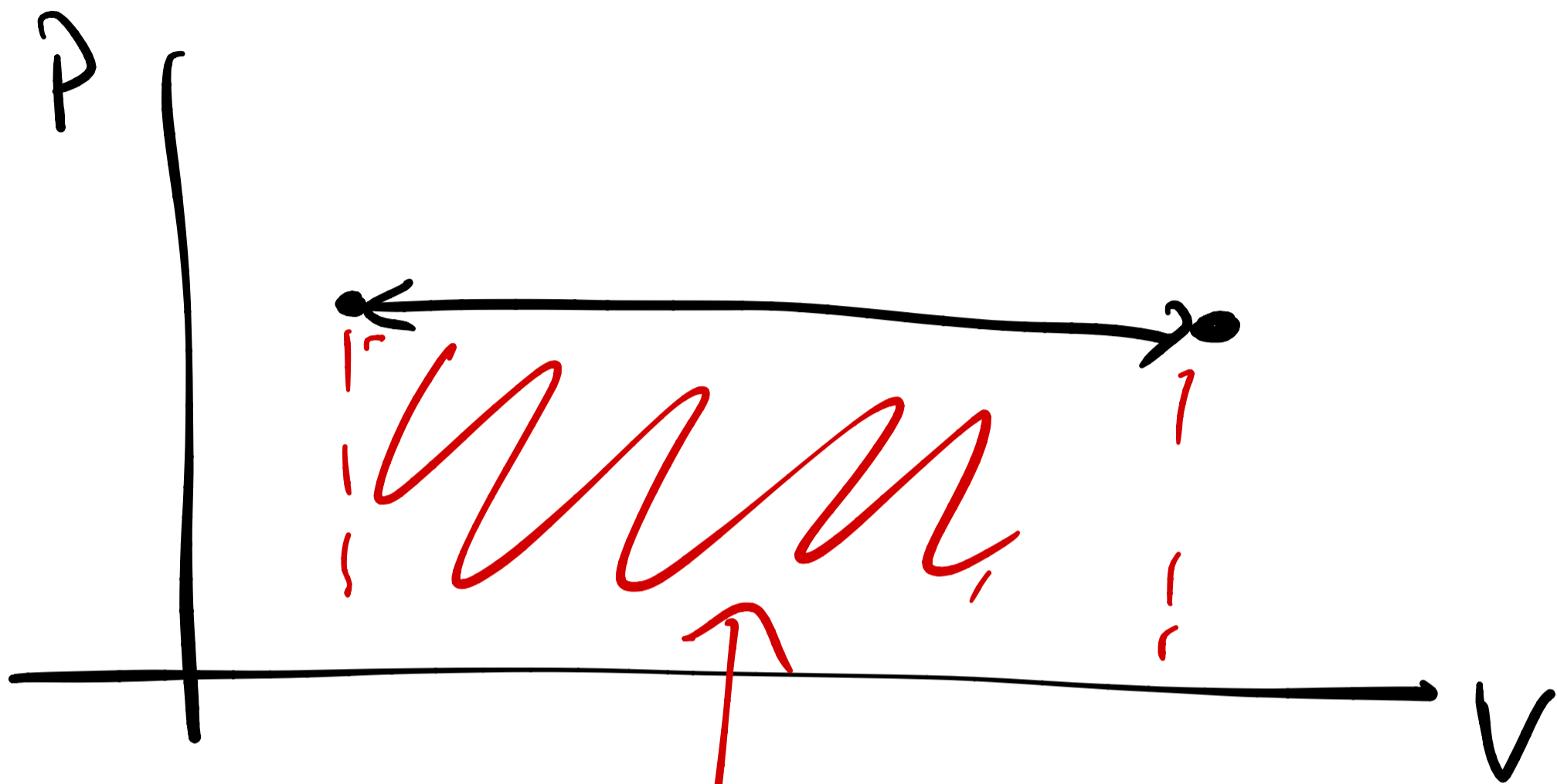
- Const T , $dT=0$

- Adiabatic, $dq=0$

Need w , q , ΔE

① dP

$$(PV = nRT)$$



true in general

$$\underline{W} = -P\Delta V$$

$$q = C_p \Delta T$$

for an ideal gas,

$$\left\{ \Delta T = \frac{P_2 V_2}{nR} - \frac{P_1 V_1}{nR} \right.$$

$$\Delta E = q + w = C_p \Delta T - P \Delta V \left. \right\} = \frac{P}{nR} \Delta V$$

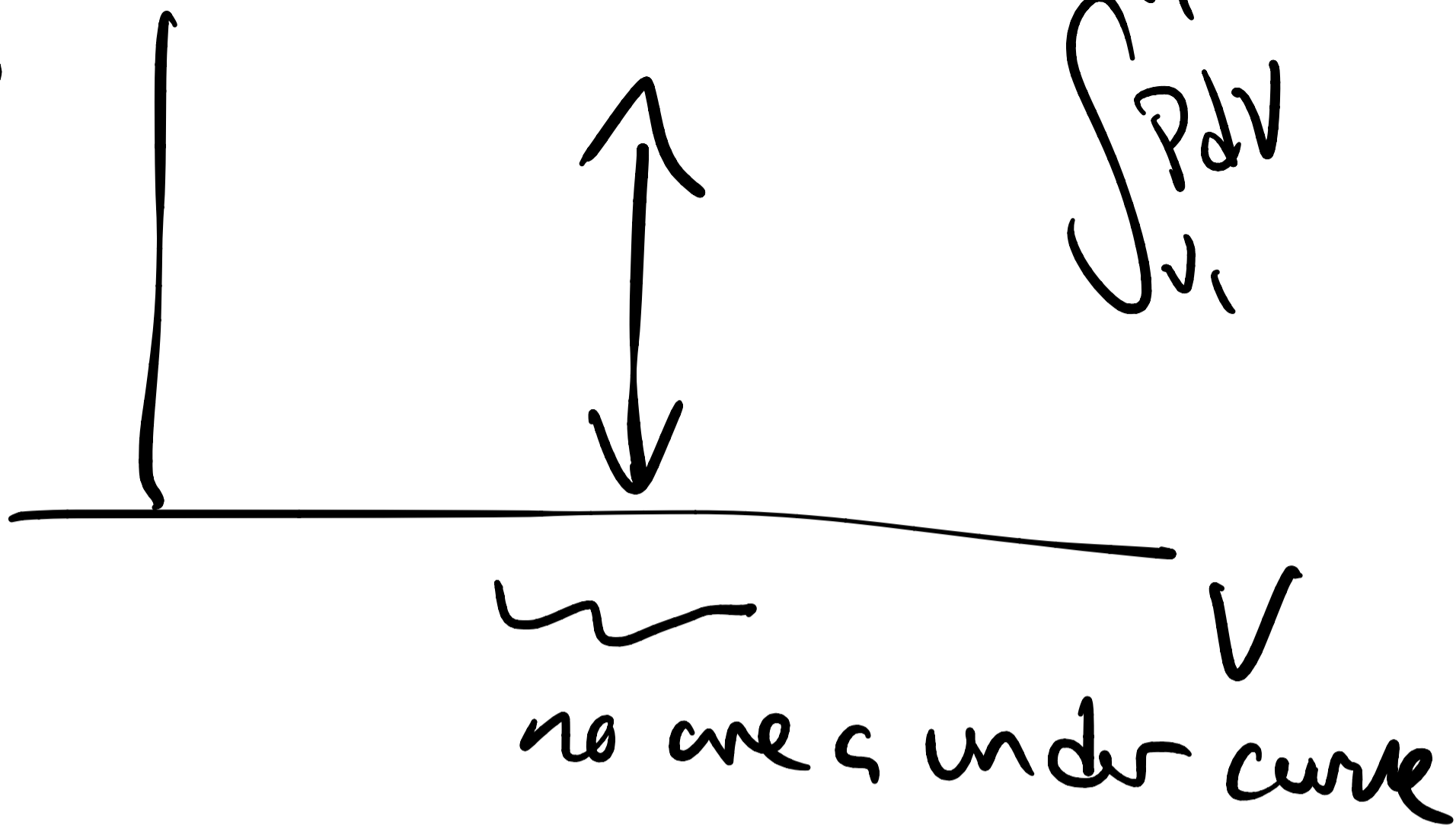
$$\textcircled{2} \quad dV = 0$$

$$(dw=0) \\ w \approx 0$$

$$q = C_V \Delta T$$

$$\Delta E = C_V \Delta T$$

q



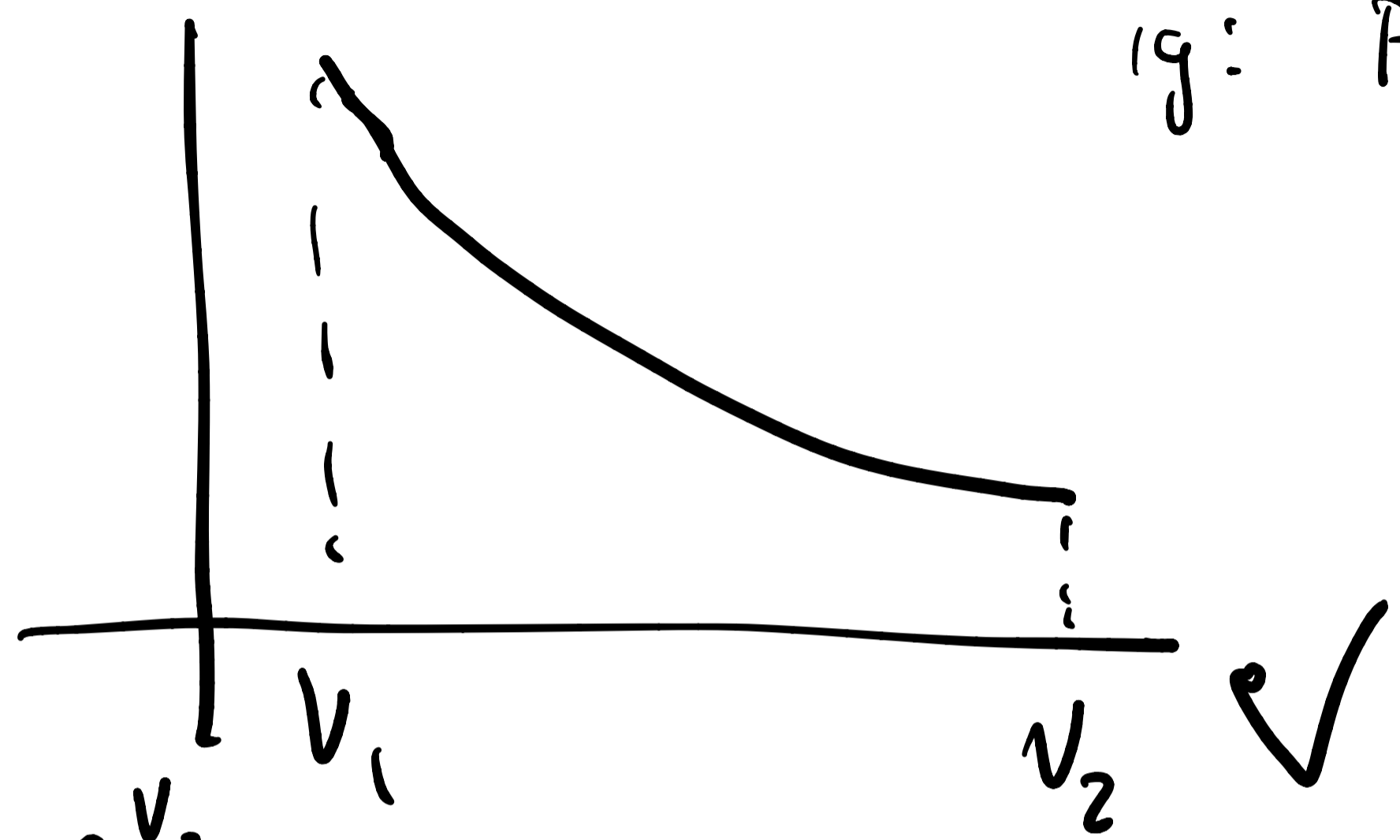
(for an ideal gas,
follows from energy
formula)

③ $dT = 0$

What does this look like

ig: $P = nRT \cdot \frac{1}{V}$

isothermal
expansion
or compression



$$\begin{aligned} w &= - \int_{V_1}^{V_2} P dV = - \int_{V_1}^{V_2} \frac{nRT}{V} dV \\ &= -nRT \int_{V_1}^{V_2} \frac{1}{V} dV \\ &= -nRT \ln(V_2/V_1) \end{aligned}$$

question
 $q > 0,$
 $q < 0,$
 $q = 0$

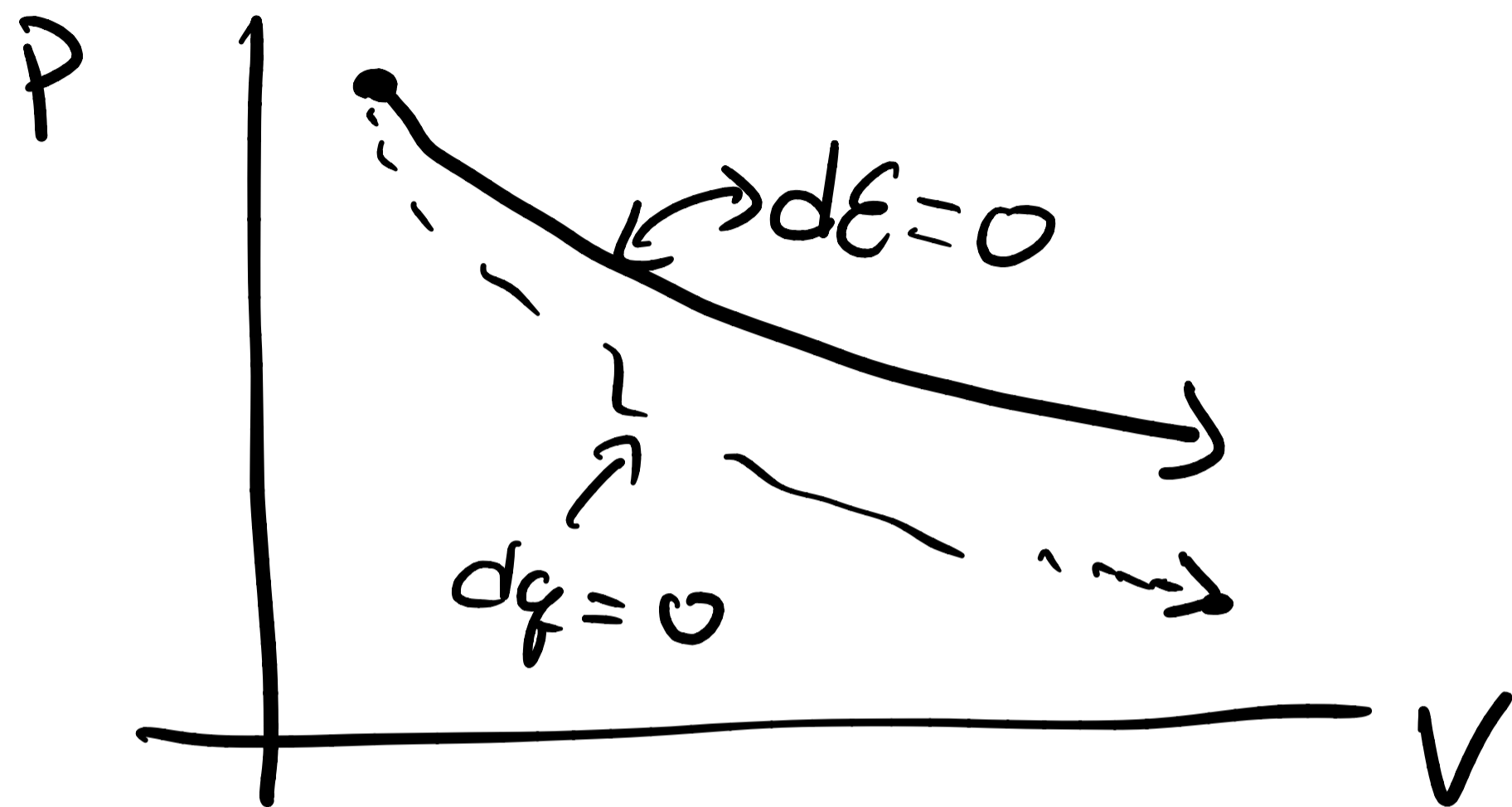
1.g. $dE = d\left(\frac{3}{2}nRT\right) = \frac{3}{2}nRdT = 0$

$$0 = dq + dw = q + w \Rightarrow q = -w$$

$$q = +nRT \ln\left(\frac{V_2}{V_1}\right)$$

$q > 0$ for expansion

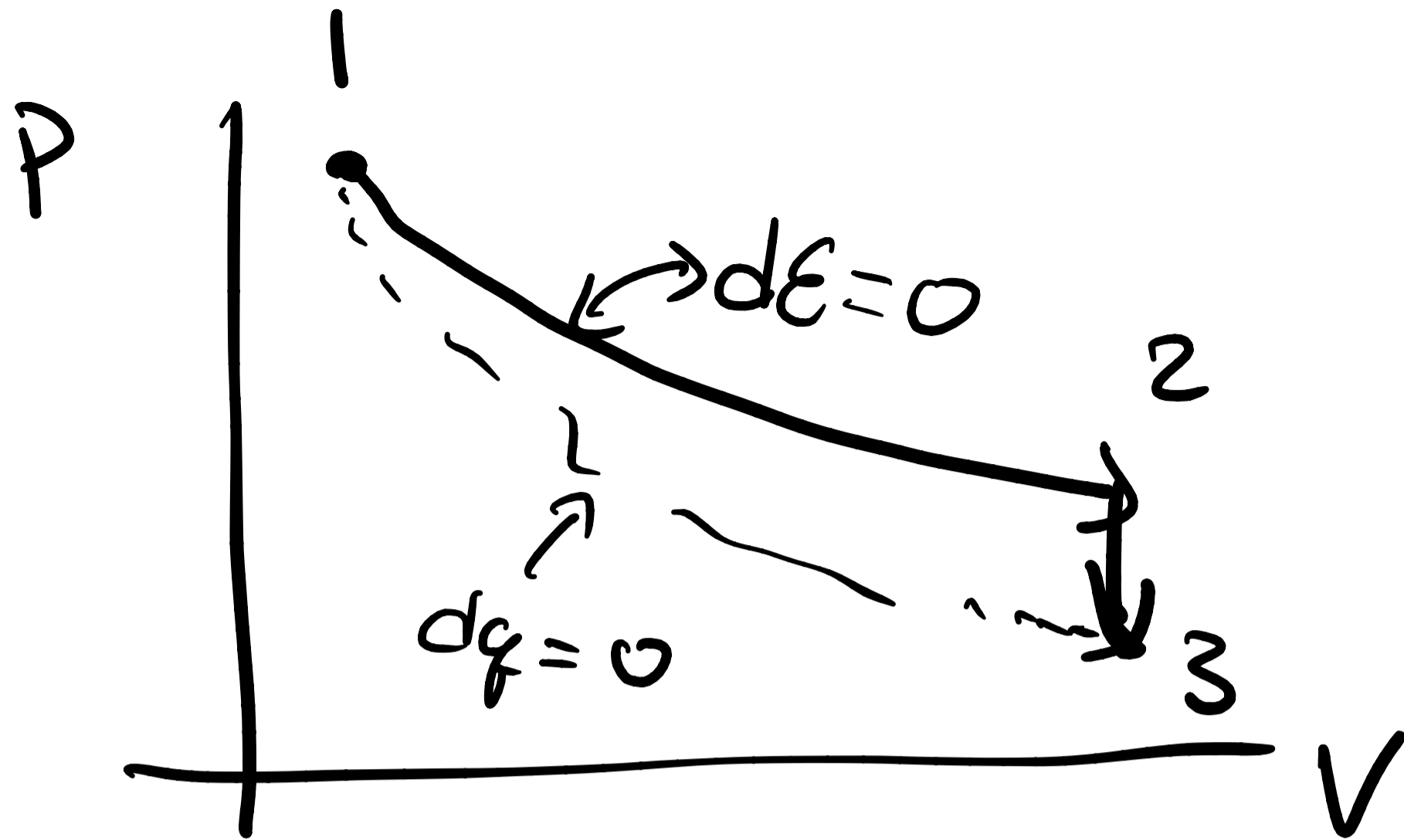
④ Adiabatic expansion, $dq = 0$



pressure
drops faster
due to cooling

$$dE = \cancel{dq} + d\omega = -P(V) dV$$
$$= -\frac{nR}{V} T(V) dV$$

how much does it cool? ← next time



$$1 \rightarrow 2 \quad dT = 0$$

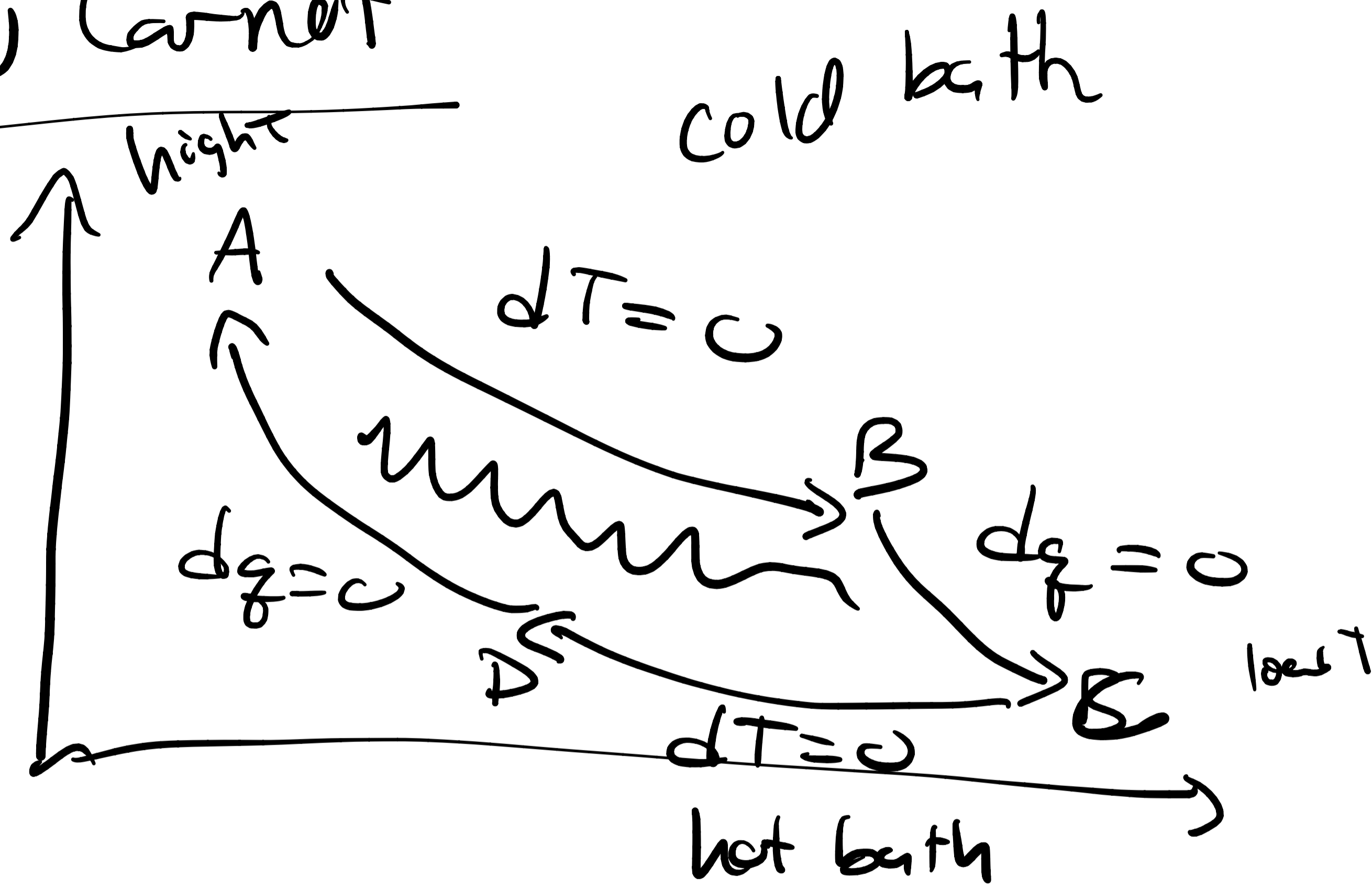
$$2 \rightarrow 3 \quad dV = 0$$

Change in E from Pt 1 \rightarrow Pt 3

$$= Pt_{1 \rightarrow 2} + Pt_{2 \rightarrow 3}$$

next time ...

Preview Carnot



4 steps

temps T_h & T_c

$$\text{efficiency} = \frac{w}{q_{in}} = 1 - \frac{T_c}{T_h}$$