

Statistical connection between
microscopic & macroscopic properties

Microscopic (atomistic) - molecular properties

x, y, z of each atom or $v_x, v_y, v_z \dots$

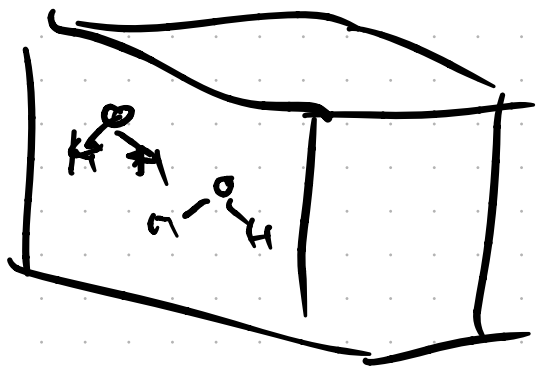
Spin of each of each atom

Macroscopic (Thermodynamic)

T, P , heat capacity, coefficient of thermal expansion

Macrostates - characterizes the thermodynamic system
 N, V, T N, p, T

Microstate - particular configuration of your system



$$N_{\text{H}_2\text{O}}, V = 1 \text{ l}$$

$$T = 25^\circ \text{C}$$

$$x_1, y_1, z_1$$

$$x_2, y_2, z_2$$

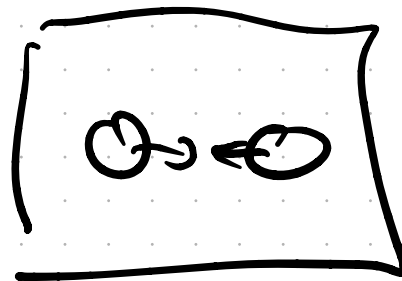
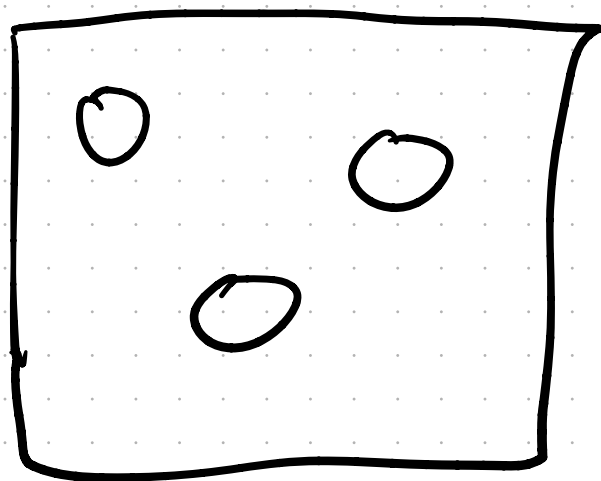
⋮

$$x_N, y_N, z_N$$

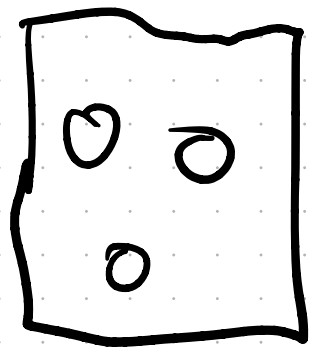
Ergodicity explores all accessible states in ∞ time

Ergodic hypothesis systems are big enough that they're effectively ergodic

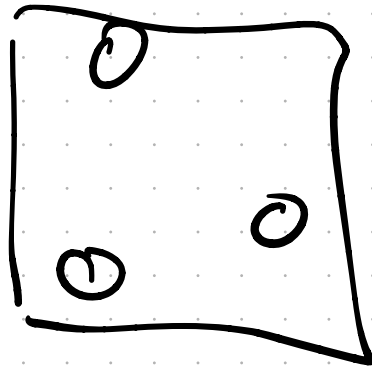
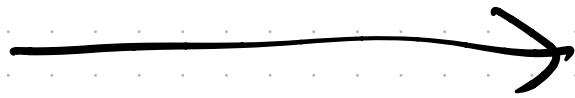
Billiard Ball Problem



Ensemble Method If you have many copies
of a system
average a property over that ensemble
Same answer as 1 system for $t \rightarrow \infty$

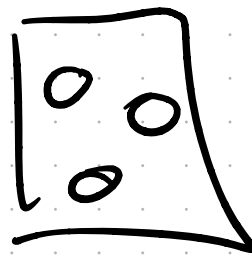
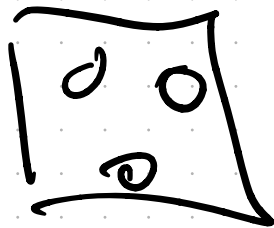
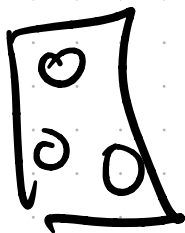
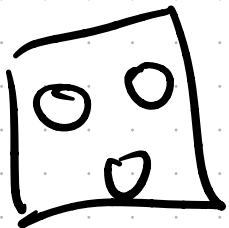


t



N, V, T

$\langle d_{12} \rangle$



ensemble

N, V, T

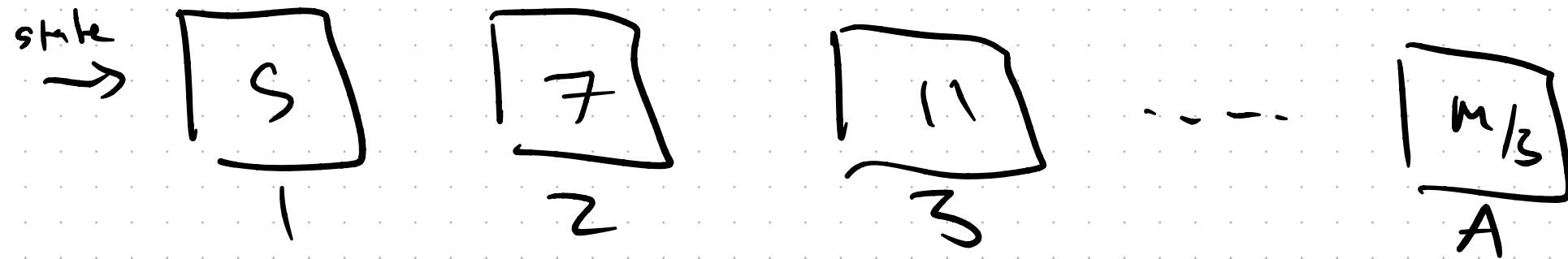
← average

have been talking about continuous systems
meaning, $\vec{x}, \vec{y}, \vec{z}$

Start with discrete systems
discrete number of possible states

Consider a system that can be
in M microstates

A total copies





N_i is the # in State i

$$\sum_{i=1}^m N_i = A$$

$$\{N_i\} = \{N_1, N_2, \dots, N_m\}$$

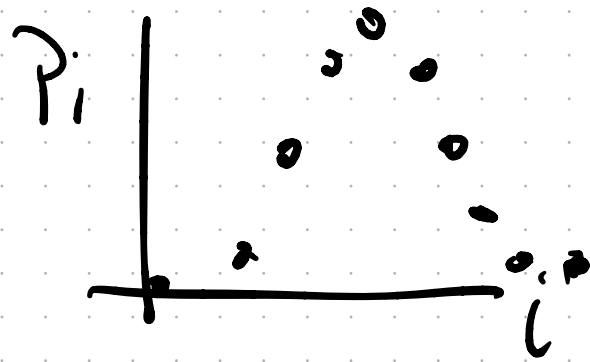
$$\{P_i\} = \{P_1, \dots, P_m\} = \left\{ \frac{N_1}{A}, \frac{N_2}{A}, \dots, \frac{N_m}{A} \right\}$$

$$m=6$$



Ensemble Average

Property y_i for each i



$$\langle y \rangle_{\text{ensemble}} = \frac{1}{M} \sum_{i=1}^M y_i P_i$$

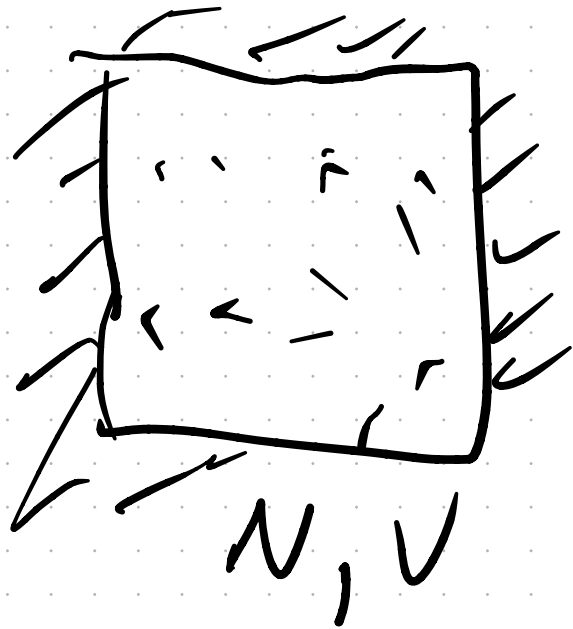
Contrast to a time average

$$y(t_1), y(t_2), \dots, y(t_{N_t})$$

$$\langle y \rangle_{\text{time}} = \frac{1}{N_t} \sum_{i=1}^{N_t} y(t_i)$$

Microcanonical Ensemble

Ensemble where every copy has the same N, U, E - isolated system



all molecules follow
Newton's equations

$$\vec{F} = m\vec{a}$$

or Schrodinger equation

Both case E is constant

Learned previously that entropy
 $S = k \ln W$ is a thermodynamic
potential for N, U, E systems

S is maximized at equilibrium
copies of this system maximize
entropy

$\Rightarrow P_1, P_2 \dots, P_m$ for our system

$$S = k_B \ln \omega$$

N copies and N_1 in state 1
 N_2 in state 2...

$$\omega = \frac{A!}{N_1! N_2! N_3! N_4! \dots N_m!} \quad (\text{multinomial distribution})$$

$$\sum_{i=1}^m N_i = A$$

A & $\{N_i\}$ large

Stirling's approximation
is okay

$$\omega = \frac{A!}{N_1! N_2! N_3! N_4! \dots N_m!}$$

$$\ln X! \approx X \ln X - X$$

$$S = k_B \ln \omega = \underbrace{(A \ln A - A)}_{-A} - \sum_{i=1}^m (N_i \ln N_i - N_i) \quad k_B$$

$$S/k_B = A \ln A - \sum_{i=1}^m N_i \ln N_i$$

$$\frac{\partial S/k_B}{\partial N_i} = 0$$

$$S \approx \frac{1}{k_B} A \ln A - \sum_{i=1}^m N_i \ln N_i$$

$$A = \sum_{i=1}^m N_i$$

$$\frac{\partial S/k_B}{\partial N_j} = A \frac{\partial \ln A}{\partial N_j} + \frac{\partial A}{\partial N_j} \ln A$$

$$\frac{\partial A}{\partial N_j} = 1$$

$$- \frac{N_j}{N_j} - 1 \cdot \ln N_j$$

$$= 1 + \ln A - 1 - \ln N_j$$

$$= \ln A - \ln N_j = \ln \left(\frac{A}{N_j} \right)$$

maximize
 N_j together
 $S \rightarrow 0$

Want to maximize all N_j

but $\sum_{j=1}^m N_j = A \leftarrow \text{constraint}$

Method of Lagrange Multipliers

$f(x) \leftarrow$ want to find max or min
with a constraint

Max
or min $\mathcal{I} = f(x) - \alpha(\text{constraint})$

$$S = k \ln \omega \quad \omega / \text{constraint}$$

$$\sum_{i=1}^M N_i = A \quad \Rightarrow \quad \left(\sum_{i=1}^M N_i - A \right) = 0$$

$$I = k \ln \omega - \alpha \left(\sum_{i=1}^M N_i - A \right)$$

$$\frac{\partial I}{\partial N_j} = \ln A - \ln N_j - \alpha = 0 \quad \text{for every } N_j$$

$$\alpha = \ln \left(\frac{A}{N_j} \right) \Leftrightarrow N_j = A e^{-\alpha}$$

$$N_i = Ae^{-\alpha}$$

$$\sum_{i=1}^m N_i - A = 0 \quad \leftarrow \text{constraint}$$

$$\sum_{i=1}^m Ae^{-\alpha} = A \Rightarrow mAe^{-\alpha} = A$$

$$e^{\alpha} = m \quad \alpha = \ln(m)$$

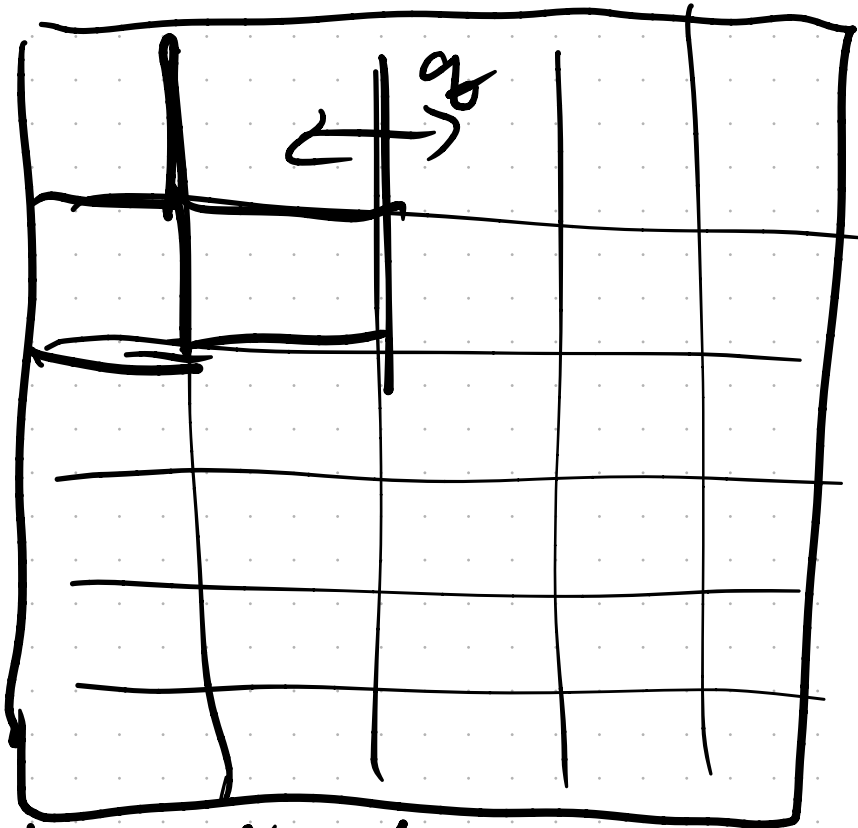
$$P_i = \frac{N_i}{A} = \frac{Ae^{-\alpha}}{A} = \frac{1}{e^{\alpha}} = \frac{1}{m}$$

Chapter 10 Canonical Ensemble

a system of interest has constant

N, U, T

isolated



N, U, E overall
heat transfers

until each
system has same T

one system @ eq

\tilde{N}, \tilde{U}, T

$$N = A\tilde{N} \quad V = A\tilde{V} \quad E = \sum_{i=1}^A \epsilon_i$$

Whole system is N, U, E

Maximize $S = k \ln \Omega$

constraints ① $\sum_{i=1}^M N_i = A$

② ~~$\sum_{i=1}^M \epsilon_i = \epsilon_{\text{total}}$~~ ϵ_{total} \leftarrow constant

Lagrange multipliers (2 constraints)

~~$S = k \ln \Omega - \alpha (\sum N_i - A) - \beta (\sum \epsilon_i - \epsilon_{\text{total}})$~~

Energy constraint

$$E_{\text{total}} = \sum_{i=1}^M \epsilon_i$$

$$\boxed{N_1}$$

1

$$\boxed{N_2}$$

2

$$E_{\text{total}} = \sum_{i=1}^M N_i \epsilon_i$$

$$S = k \ln \omega - \alpha (\sum N_i - A)$$

$$- \beta (\sum N_i \epsilon_i - E_{\text{total}})$$

$$\frac{\partial S}{\partial N_j} = \ln(A/N_j) - \alpha - \beta \epsilon_j = 0$$

$$\ln(A/N_j) - \alpha - \beta \epsilon_j = 0$$

$$N_j = e^{-\alpha} e^{-\beta \epsilon_j} \cdot A$$

$$\sum_{j=1}^m N_j = A = A e^{-\alpha} \sum_{j=1}^m e^{-\beta \epsilon_j}$$

$$1 = e^{-\alpha} \sum_{j=1}^m e^{-\beta \epsilon_j}$$

$$e^{\alpha} = \sum_{j=1}^m e^{-\beta \epsilon_j}$$

$$P_j = \frac{N_j}{A} = e^{-\alpha} e^{-\beta \epsilon_j}$$

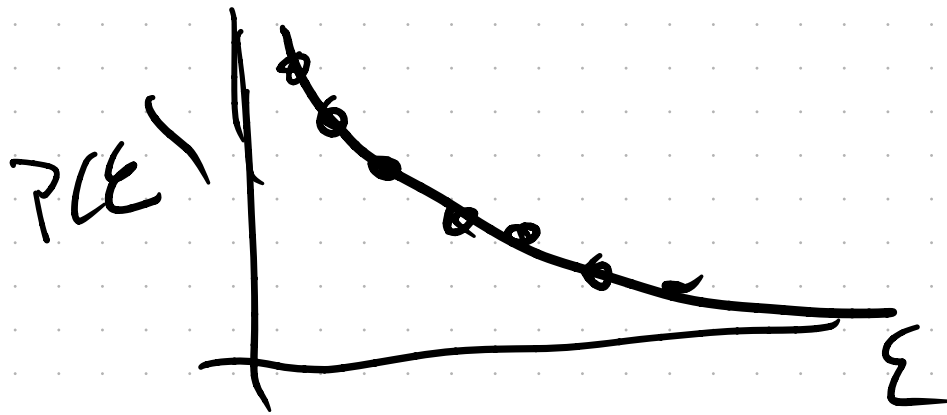
$$= \frac{e^{-\beta \epsilon_j}}{\sum_{j=1}^m e^{-\beta \epsilon_j}}$$

$$P_j = \frac{e^{-\beta \epsilon_j}}{\sum_i e^{-\beta \epsilon_i}}$$



$$Q = \sum_i e^{-\beta \epsilon_i}$$

"partition function"



$e^{-\beta \epsilon_j}$ is called
Boltzmann factor
(relative weight)

for Micro Canonical

Boltzmann factor = 1 like $\beta = 0$

$$Q = \sum_{i=1}^M \beta F_i \approx M$$

$$\langle y \rangle = \sum_{i=1}^M y_i P_i = \frac{\sum_{i=1}^M y_i e^{-\beta E_i}}{Q}$$

$$\langle E \rangle = \frac{\sum_{i=1}^M E_i e^{-\beta E_i}}{Q}$$

$$\beta = 1/k_B T$$

$$Q = \sum_{i=1}^M e^{-\beta \epsilon_i}$$

can derive thermodynamic quantities
from the partition function

$$\langle \epsilon \rangle = \frac{\sum_{i=1}^M \epsilon_i e^{-\beta \epsilon_i}}{Q}$$

$$-\frac{\partial Q}{\partial \beta} = \sum_{i=1}^M \epsilon_i e^{-\beta \epsilon_i}$$

$$\frac{\partial \log Q}{\partial x} = \frac{1}{Q} \frac{\partial Q}{\partial x}$$

$$-\frac{\partial \ln Q}{\partial \beta} = \frac{\sum_i \epsilon_i e^{-\beta \epsilon_i}}{Q} = \langle \epsilon \rangle$$

$$\langle \epsilon \rangle = -\frac{\partial \ln Q}{\partial \beta}$$

Gibbs Entropy

$$S = k \ln \Omega_{\text{micro}}$$

$$S = -k_B \sum_{i=1}^M p_i \ln p_i \quad \sum p_i y_i = \langle y \rangle$$

$$p_i = e^{-\beta \epsilon_i} / \Omega \quad \nearrow$$

$$= -k_B \sum p_i [-\beta \epsilon_i - \ln \Omega]$$

$$= k_B \beta \sum p_i \epsilon_i + k_B \ln \Omega$$

$$= k_B \beta \langle \epsilon \rangle + k_B \ln \Omega \quad \leftarrow \text{entropy}$$

$$S = k_B \beta \langle \epsilon \rangle + k_B \ln Q$$

$$T = \left(\frac{\partial U}{\partial S} \right)_V = \frac{\partial U}{\partial S}$$