Lecture 7 - Intro to smpling Reminder: goal of stat merch Is to compute average observables Loyensmille = Jdx O(X) P(X) If const NVT, $P(x) = e^{-\beta H(x)}/2$ Some problems can solve exactly, put most not. Will talk about how to solve numerically First: finish canonical ensuble $A = -k_BT \ln 2$ Also knew P= - 2A/2V ctc 19-16 -= 3 40

first flacsh ideal gas: $Q = \frac{1}{N!} \left(\frac{V}{\Lambda^3} \right)^N = \frac{1}{N!} \left(\frac{V}{2\pi m k_g T} \right)^{\frac{N}{2}N}$ $= \beta^{-3/2} N \cdot 5 + \gamma P P$ E = - ShQ = ZN ShP = ZN kgT V one more thomas quartity of Importance! $C = \frac{\partial \mathcal{E}}{\partial T} \qquad C_{\nu} = \begin{pmatrix} \partial \mathcal{E} \\ \partial \mathcal{F} \end{pmatrix}_{\nu} \qquad C_{\rho} = \begin{pmatrix} \partial \mathcal{E} \\ \partial \mathcal{F} \end{pmatrix}_{\rho}$ Idenlgas! Cu= Zuke In general, C = DE = DE DF = The second sec = -1 22 HATZ JP = -KB P 22E

Now, for NVT bor E= - 2 lat Op $C_{v} = -k_{b} \beta^{2} \frac{\partial \varepsilon}{\partial \beta}$ $= k_{B} \beta^{2} \frac{\partial}{\partial \beta} \left[\frac{1}{2} \frac{\partial}{\partial \beta} \right]$ $= k_{0} p^{2} \left[-\frac{1}{2} \frac{\partial }{\partial p} \frac{\partial }{\partial p} + \frac{1}{2} \frac{\partial ^{2} }{\partial p} \right]$ $= k B B^{2} [\langle E^{2} \rangle - \langle E \rangle^{2}]$ = KBB2 Var(E) $Z = \int dx e^{-\beta \mathcal{H}(x)} \frac{32}{5\beta^2} = \int dx \mathcal{H}(x) e^{-\beta \mathcal{H}(x)}$ DE = L Vor(E) &

Example of floctuction dissipation theorem: fluctuations @ cq prop to how quantity recursitively (Onsuger Negnession) What does vor (E) men physically? ECA MAJANA Joe Energy goves in Bart of bath How big is floc relative to E? JE/E = JEE X JEV/E & JI P/C COAN

These energy fluctuations are what allow cherical systems to overcome barners and endryd reactions. But how do ve calculate properties if we can't explicitly get 7? Consider the problem of computing く0>=」(Jax OCA) PCA) first, lets consider that most Observables may only deput on pisition $\langle 0 \rangle = \int J x^{3} f p^{3N} O(\vec{x}) c / \frac{1}{2}$ and $\mathcal{H} = \Sigma f^{3} (m + U(\vec{x}))$

Then $\langle 0 \rangle = const \cdot \int d\vec{x} e^{-\beta n c \vec{x}} O(\vec{x})$ Candeline cnother pos. Z= Jdxe fucri $\langle 0 \rangle = \int d\vec{x} \, \partial \vec{x} \, e^{-\mu x} / \gamma$ $Z = cmst \tilde{z}$ Nou consider 12 probles (an compute Co> numerically "by guadrahme" ie イロンニ を O(x;) P(x;) &×

Seems great, what the problem
For each dimension:
points =
$$\frac{1}{\Delta x}$$

So Por & dimensions
= $(\frac{1}{\Delta x})^d = c^d \ln(\frac{1}{\Delta x})$
= $(\frac{1}{\Delta x})^d = c^d \ln(\frac{1}{\Delta x})$
expansibly large $(\frac{d}{d} = 3N)$ gets
big first)
I dee of sampling:
Generate representative effs
st $\chi_{\ell} \propto P(\chi_{\ell})$ Sorehow
Thun $\mathcal{L}O \gg \frac{1}{T} \underbrace{\mathcal{L}}O(\chi_{\ell})$
One way to do this is "Molecular
Dynamics", minic Newtont tenf

First discuss Monte Carlo guvert a "Markon Chain" Xf -> Xt+1, depends only on connert State If this satisfies "detailed blue" $P(x_{t})P(x_{t} \rightarrow x_{t+1}) = P(x_{t+1})P(x_{t} \rightarrow x_{t})$ flun con be proven that Chain Counzes s.t. X~P(x)