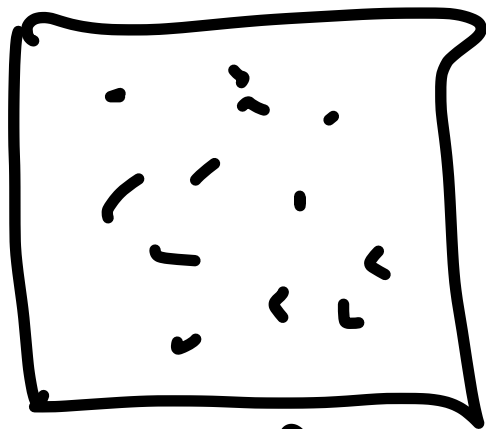


# Lecture 4 - Thermo Review

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Last time:

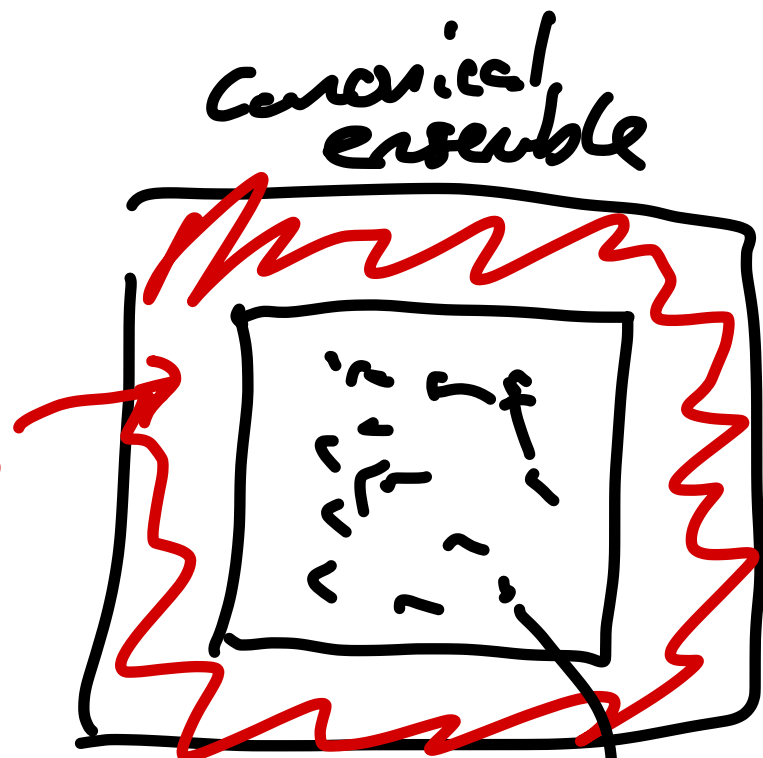
- Microcanonical partition function



$N, V, E$   
constant

closed  
isolated

bath  
temp  
 $= T$



heat flows

$N, V, T$  constant

- Asserted that (microcanonical)

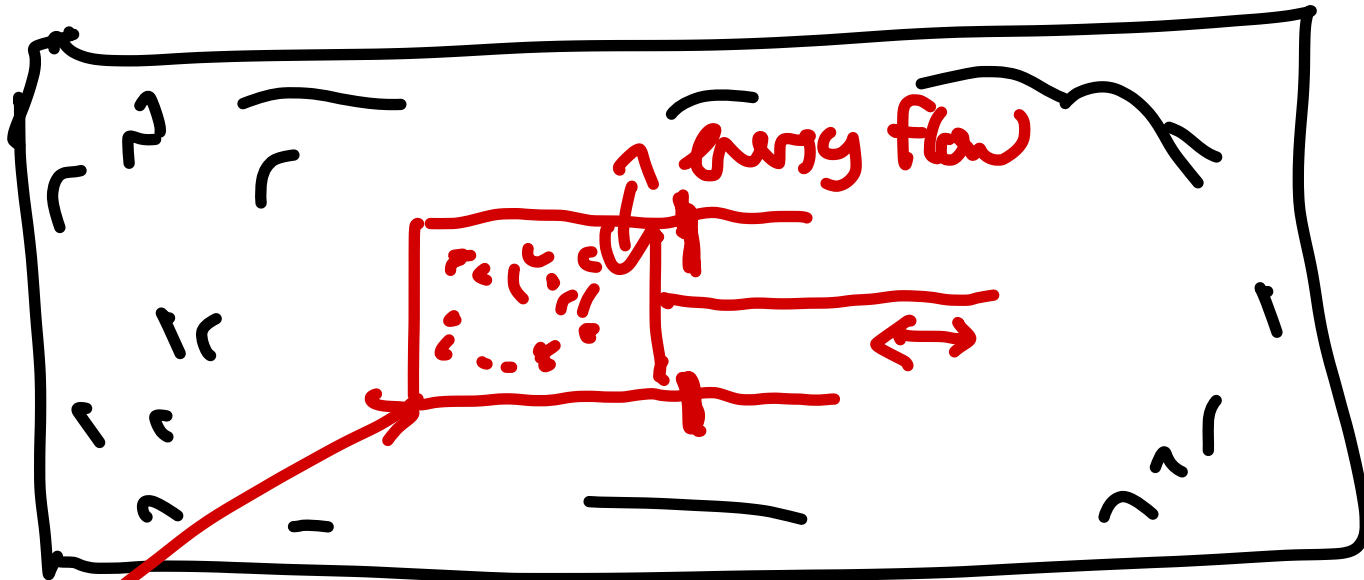
$$S = k_B \ln \Omega(N, U, E)$$

## Thermo review

- ① First law: conservation of energy. Energy isn't created or destroyed

What is the system we're studying

Classic scenario:



closed, isolated  $N, V, E$  constant

System

const  $N, V, T$

we care about

System - does not have const  $E$

another Statement:

for a system

change  $E_{\text{system}} =$  heat flows in  
- work done  
by system

Change  $E_{\text{sur}} = -$  Change  $E_{\text{system}}$

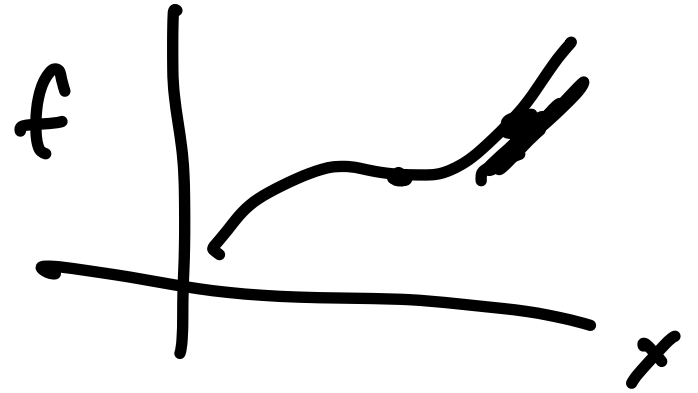
$$dE_{\text{system}} = \delta q - \delta w$$

↑ heat in                      ↑ work done

Reminder: "infinitesimal" ( $\epsilon, d\epsilon$ )

response to a tiny change in  
of a function argument

$f(x)$  is a function  
 $df$



$$f(x + \Delta x) \approx f(x) + \Delta x \frac{df}{dx} + \frac{1}{2} \Delta x^2 \frac{d^2 f}{dx^2} \dots$$

$$df = \lim_{\Delta x \rightarrow dx} f(x + \Delta x) - f(x) = \left( \frac{df}{dx} \right) dx$$

$f(x, y)$  what is  $df$

$$f(x + \Delta x, y + \Delta y) = f(x) + \Delta x \left( \frac{\partial f}{\partial x} \right)_y + \Delta y \left( \frac{\partial f}{\partial y} \right)_x + \dots$$

$$df = \left( \frac{\partial f}{\partial x} \right)_y dx + \left( \frac{\partial f}{\partial y} \right)_x dy + \dots$$

connected to chain rule, eg

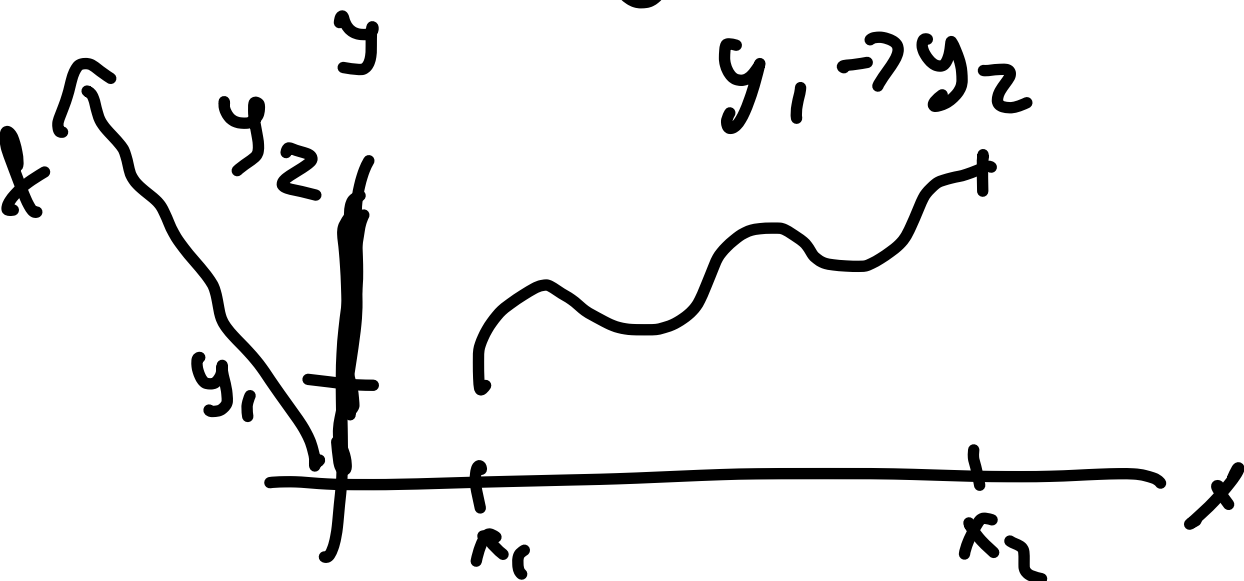
$$\frac{df(x, y)}{dt} = \left( \frac{\partial f}{\partial x} \right) \left( \frac{\partial x}{\partial t} \right) + \left( \frac{\partial f}{\partial y} \right) \left( \frac{\partial y}{\partial t} \right)$$

Thermodynamics - want to know  
change in  $f$  along a thermodynamic  
path

$$\Delta f = f(x_2, y_2) - f(x_1, y_1)$$

$$= \int_{\substack{x_1 \rightarrow x_2 \\ y_1 \rightarrow y_2}} df$$

(eg  
 $\mathcal{E}(V, T)$ )



$$dE = \delta q - \delta w$$

Energy is a "state variable"  
heat, work are not

$$\Delta E = \int_{\text{path}} \delta q - \delta w$$

different kinds of work  $\left(\frac{\partial u}{\partial r}\right)$   
mechanical, chemical

classical  
mech.

$$W = \int_{\text{path}} F \cdot dr$$

$$dw = F \cdot dr$$

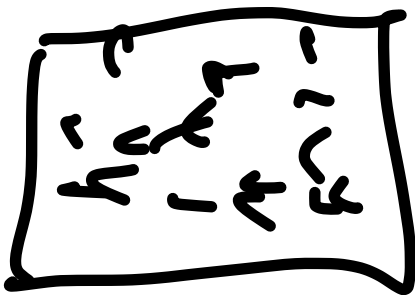


$$d\omega = - \left( \frac{\partial u}{\partial r} \right) dr \quad \text{energy}$$

$$E(A, B, C) \quad \left( \text{think } A, B, C = N, U, T \right)$$

$$dE = \left( \frac{\partial E}{\partial A} \right)_{B, C} dA + \left( \frac{\partial E}{\partial B} \right)_{A, C} dB + \left( \frac{\partial E}{\partial C} \right)_{A, B} dC$$

1st law  $dE = \delta q - \delta w$  ↙ always



$$\Rightarrow -\delta w = \sum_i \left( \frac{\partial E}{\partial \lambda_i} \right) d\lambda_i$$

closed isolated system -  $\delta q = 0$   
 ' intensive

Real example

$$-w = -PdV + \mu dN$$

chemical work

② 2nd law:

heat is not a state function

exists a quantity  $dS = \frac{\delta Q}{T}$

• state function, "entropy"

$$S(b) - S(a) = \int_{a \rightarrow b} \frac{\delta Q}{T}$$

units of  
 $[E]/[T]$

$$d\varepsilon = \delta q - \delta w$$

$$= \delta q + \sum_i \left( \frac{\partial \varepsilon}{\partial \lambda_i} \right) d\lambda_i$$

divide  
by T

$$\frac{1}{T} d\varepsilon = \underbrace{\frac{\delta q}{T}}_{dS} + \frac{1}{T} \sum_i \left( \frac{\partial \varepsilon}{\partial \lambda_i} \right) d\lambda_i$$

$$dS = \frac{1}{T} d\varepsilon - \frac{1}{T} \sum_i \left( \frac{\partial \varepsilon}{\partial \lambda_i} \right) d\lambda_i$$

$$dS = \frac{1}{T} dE - \frac{1}{T} \sum_i \left( \frac{\partial E}{\partial \lambda_i} \right) d\lambda_i$$

in microcanonical ensemble

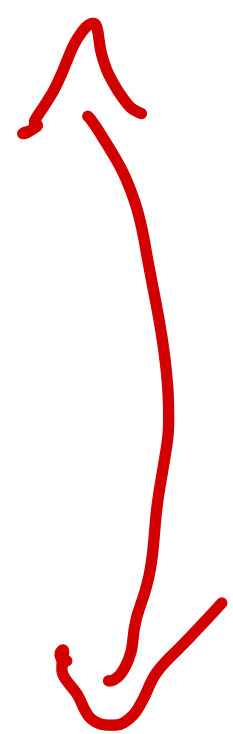
$S$  depends only on  $N, V, E$

infinitesimals - chain rule ..

$$S(N, V, E) \Rightarrow$$

$$dS = \left( \frac{\partial S}{\partial N} \right)_{V, E} dN + \left( \frac{\partial S}{\partial V} \right)_{N, E} dV + \left( \frac{\partial S}{\partial E} \right)_{N, V} dE$$

$$\Rightarrow \left( \frac{\partial S}{\partial E} \right)_{N, V} = \frac{1}{T}$$



$$dS = \left( \frac{\partial S}{\partial N} \right)_{V, E} dN + \left( \frac{\partial S}{\partial V} \right)_{N, E} dV + \left( \frac{\partial S}{\partial E} \right)_{N, V} dE$$

$$dS = \frac{1}{T} dE - \frac{1}{T} \left( \frac{\partial E}{\partial N} \right)_{V, E} dN - \frac{1}{T} \left( \frac{\partial E}{\partial V} \right)_{N, E} dV$$

$$\Rightarrow \left( \frac{\partial S}{\partial N} \right)_{V, E} = \frac{\mu}{T}$$

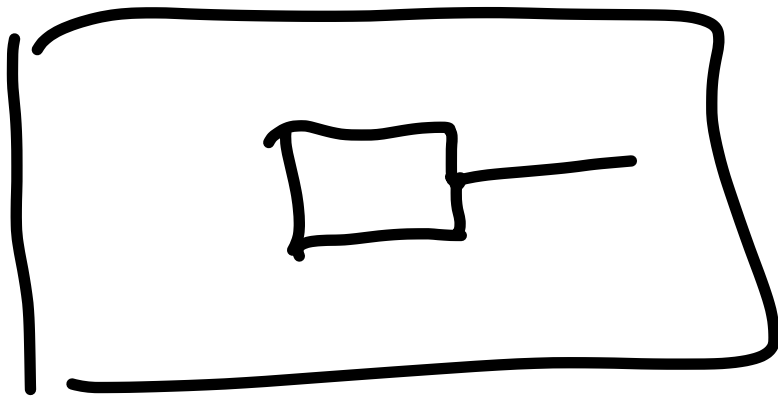
$$\left( \frac{\partial S}{\partial V} \right)_{N, E} = \frac{P}{T}$$

•  $\Delta S \geq 0$  for any process

$\Delta S = 0$  for a reversible path

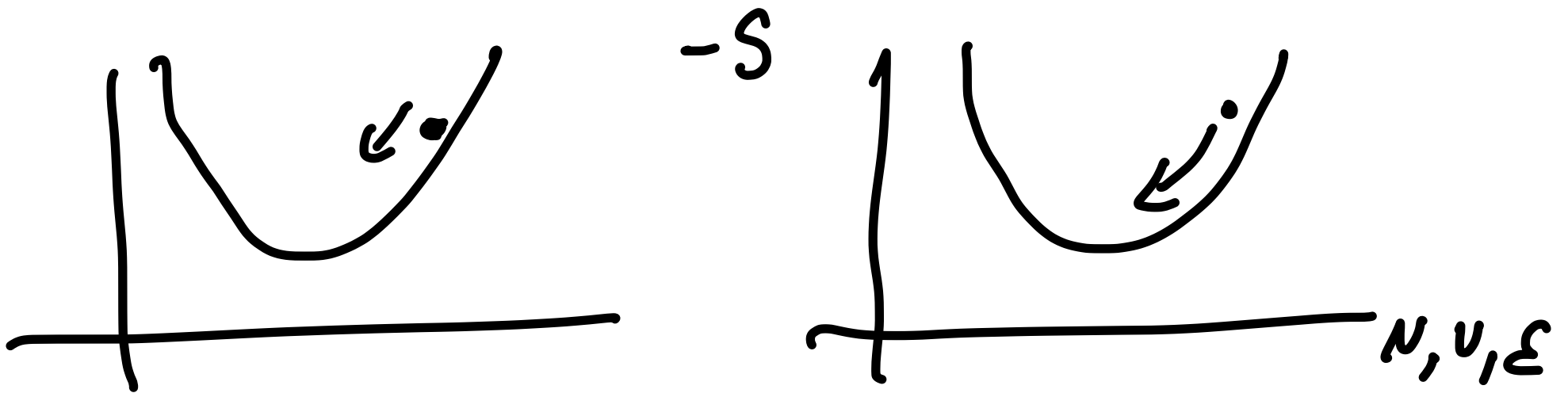
includes isolated systems

universe = isolated system  
constant energy



Isolated system

Can think of  $S$  is  $[E-S]$   
a "thermodynamic potential"  
for microcanonical ensemble



entropy will maximize as  
we approach equilibrium

Next time:  $S$  for ideal gas  
ideal gas law from  $S \dots$