

Lecture 19

Important quantities in phase transitions
are correlation functions of quantities in space

Eg $g(r)$ - if particle is at $r=0$, what is
(relative) likelihood of seeing one at distance r

We found we could write this as

$$\langle \rho(0) \rho(\vec{r}) \rangle \quad \text{where} \quad \rho(\vec{r}) = \sum_{i=1}^N \delta(\vec{r} - \vec{r}_i)$$

In a lattice model, compute spin-spin correlation
function - \uparrow - likely to be \uparrow nearby but how far?

$$\text{Let } \delta s_i = s_i - \langle s_i \rangle$$

$$C_{ij} = \langle \delta s_i \delta s_j \rangle = \langle s_i s_j \rangle - \langle s_i \rangle \langle s_j \rangle$$

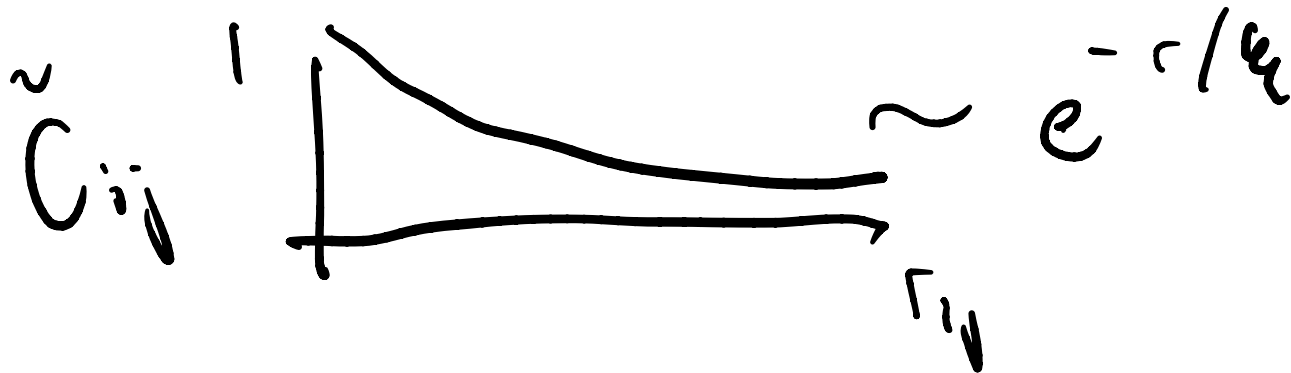
if i & j very far apart, could be independent

$$\text{So then } \langle s_i s_j \rangle \rightarrow \langle s_i \rangle \langle s_j \rangle$$

So $C_{ij} \rightarrow 0$ in this case

Note $C_{ii} = \text{Var}(s_i)$ so define

$$\tilde{C}_{ij} = C_{ij} / \text{Var}(s_i)$$



Volume
of d correlated

spins correlated w/ $i \approx \sum_{j \neq i} C_{ij}$

This is related to susceptibility in mag:

$$\chi = \frac{1}{N} \text{Var}(\delta M) \quad \delta M = \sum_{i=1}^N (s_i - \langle s_i \rangle)$$

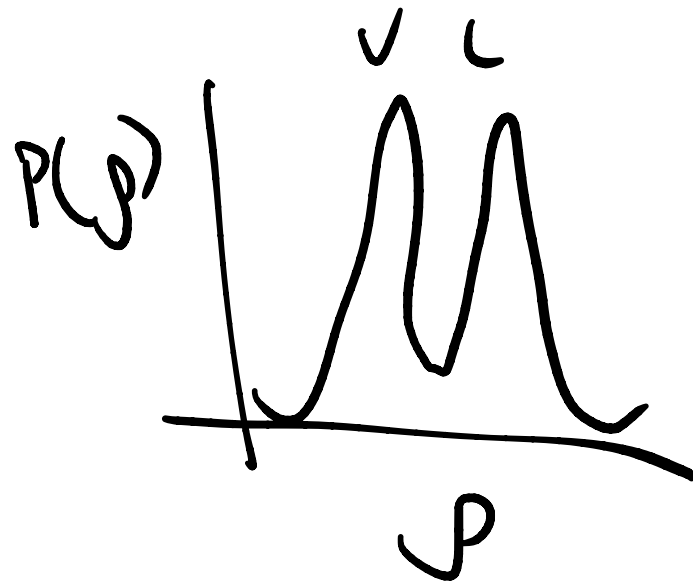
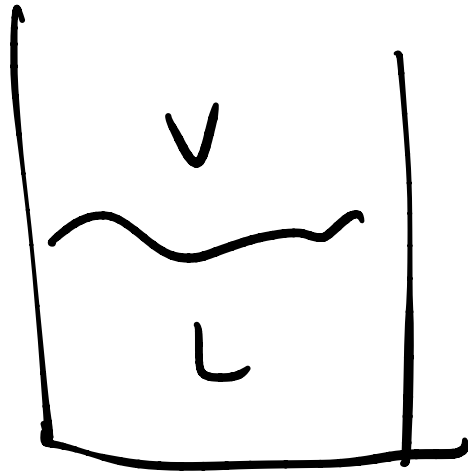
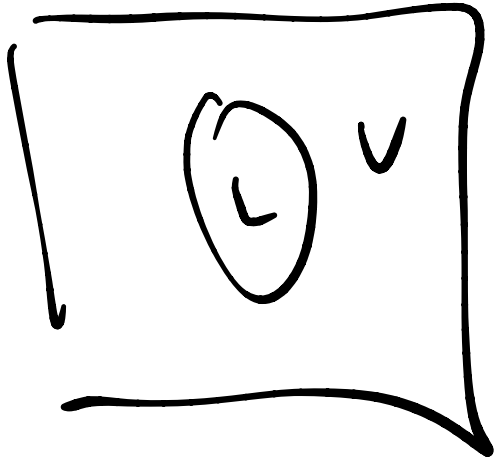
$$= \frac{1}{N} \sum_{i,j} \langle s_i s_j \rangle - \langle s_i \rangle \langle s_j \rangle$$

$$= \frac{1}{N} \sum_i \sum_j \langle s_i s_j \rangle - \langle s_i \rangle \langle s_j \rangle = \sum_j C_{ij}$$

Since χ diverges at phase transition, means correlation volume grows

2 ways V or L can diverge

1) 2 phase coexistence



as $N \rightarrow \infty$, PLP gets sharp & V or L diverges

2) cts phase transition: no distinction,

divergence means long range spatial correlations

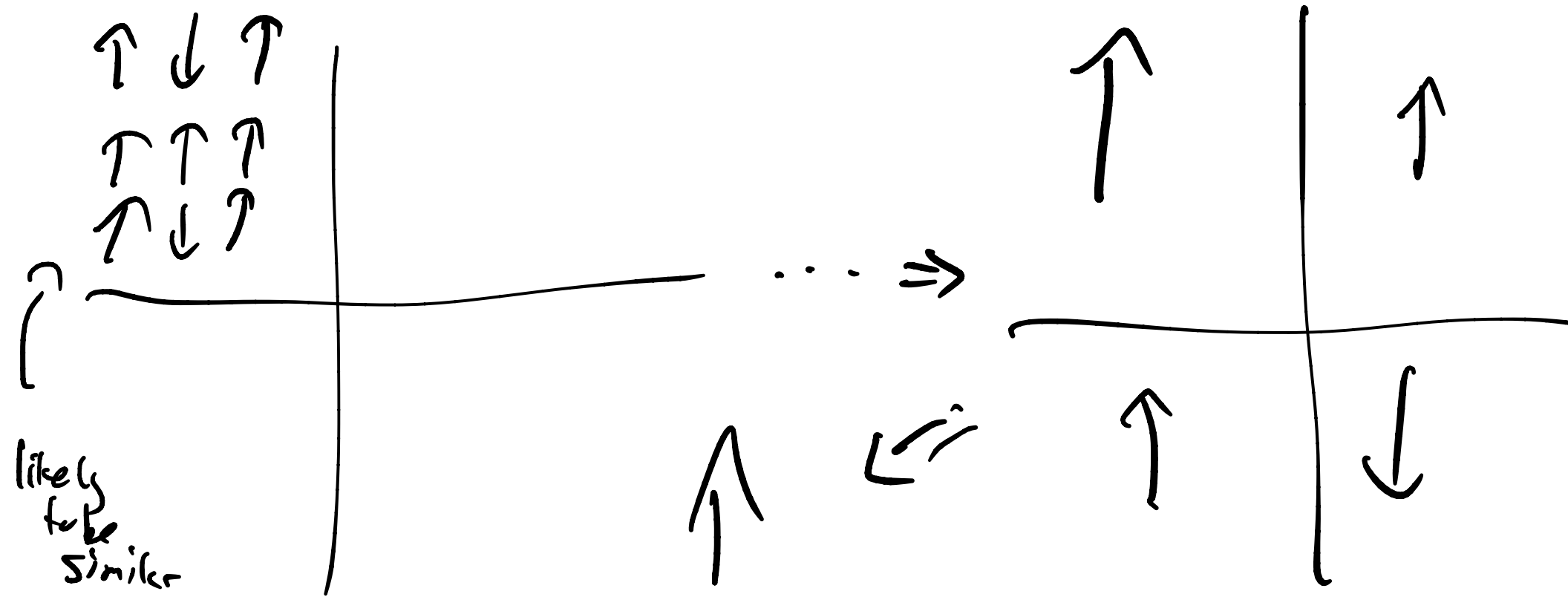
fit to $G(r) \sim C_{ij} \sim e^{-r/\xi} / r^{d-2+\zeta}$

as $T \rightarrow T_c$ $\xi \sim |T - T_c|^{-\nu}$

Large length scale \Leftrightarrow system looks similar
on large & small length scales

Renormalization Group

coarse grain over part of system to
get partition function that looks same



Eg block spin $H = \sum_{\langle i,j \rangle} J s_i s_j \rightarrow \sum_{\langle i',j' \rangle} J' s_{i'} s_{j'}$

Repart: if $J \rightarrow J'$ process converge,
 then there is a fixed point, & this corresponds
 to a phase transition

From eqns that generate $J \rightarrow J'$,
can compute critical exponents!

For 2d ising, 0 field, find two stable
fixed points ($T=0, T=\infty$)
& 1 unstable fixed point

$$J/k_B T_c = 0.50698 \quad (\text{Chandler Ch 5})$$

where exact is 0.44069 ↙ Book 261

iterations needed gets larger as closer to T_c ,
connecting to a growing length scale

Intro to non-eg

So far we dealt with systems after they reach equilibrium

eg heat flows in & out until $T_{\text{bath}} = T_{\text{sys}} = T$,
doing work on the system

What happens during this process?

- time reversibility broken (at macro level)
- can reach non eg steady state
w/ const driving but basically like eg

Eg! self-assembly by drying
in external field

molecular motors (ATP gradient)
folding/unfolding protein under force

Some theories for these real cases

First here to understand near- eq & dynamics
time dep processes

Brownian motion & Langevin eqn

Observation: pollen particles jiggled randomly
predictable behavior on average (Brown)

Einstein postulated, maybe due to
random collisions w/ solvent

$$F_{\text{total}} = m a = m \frac{dv}{dt}$$

physical observations / Stokes - low drag

$$m \frac{dv}{dt} = -\zeta v$$

$$\zeta = 6\pi \eta a$$

radius
for sphere