

Phase transitions :

Part 2

Last time: discussed general physics of phase transitions

We introduced the Ising Model as a simple way to study magnetization h.c. transition

$$H = -\frac{1}{2} \sum_{\langle i,j \rangle} JS_i S_j - h \sum_i S_i \quad (\text{n dimensions})$$

$$\propto 1d, \omega / PBC \quad H = - \sum_{i=1}^N (JS_i S_{i+1} + \frac{h}{2}(S_i + S_{i+1}))$$

Q: When is there a spontaneous magnetization transition: $m = M/N = \langle \sum S_i \rangle / N = \langle S_i \rangle > 0, \text{ or } h = 0$

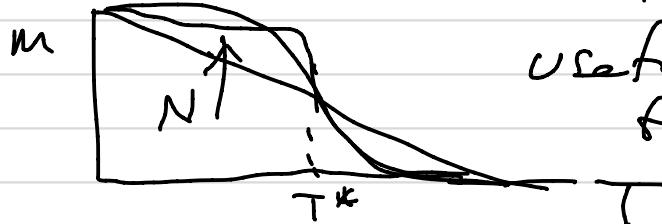
Said, should not exist for 1d ising model

For 1d model, $h=0$, said

$$Z \propto (2 \cosh(\beta J))^N$$

$$\text{and } f = F/N = -k_B T/N \log Z = -\frac{1}{\beta} \log(2 \cosh(\beta J))$$

Note: w/ $h > 0$



useful model /
for binding, e.g.

But what about $h \neq 0$?

Need a new technique called transfer matrices

$P_{S,S'}$ is the matrix w/ entries $e^{\beta j(Ss') + \beta h(S+s')/2}$

$$\text{so } \langle 1|P|1\rangle = e^{\beta j + \beta h} \quad \langle -1|P|-1\rangle = e^{+\beta j - \beta h}$$

$$\langle 1|P|0\rangle = \langle 0|P|1\rangle = e^{-\beta j}$$

$$P = \begin{pmatrix} e^{\beta(j+h)} & e^{-\beta j} \\ e^{-\beta j} & e^{\beta(j-h)} \end{pmatrix}$$

$$\text{so } Z = \sum_{S_1, \dots, S_N} \langle S_1 | J | S_2 \rangle \langle S_2 | J | S_3 \rangle \dots \langle S_N | J | S_1 \rangle$$

Spin eigenvectors are complete ie $\sum_{S_i} |S_i\rangle \langle S_i| = 1$

$$\Rightarrow Z = \sum_{S_1} \langle S_1 | P^N | S_1 \rangle = \text{Tr}(P^N)$$

$\text{Tr}(M^n) = \text{Tr}(UD^nU^{-1}) = \sum_{i=1}^N \lambda_i^n$, λ_i are eigenvalues of M if n diagonalizable

$$\Rightarrow Z = \lambda_1^n + \lambda_2^n \quad \text{Det}[P - \lambda I] = 0$$

$$\lambda = e^{\beta j} (\cos(\beta h) \pm \sqrt{\sinh^2(\beta h) + e^{-4\beta j}})$$

$$\text{as } N \rightarrow \infty \quad Z \approx \lambda_+^n$$

$$\sinh(x) = \frac{e^x - e^{-x}}{2} \quad \cosh(x) = \frac{e^x + e^{-x}}{2}$$

and $f(h, \beta) = -\frac{1}{\beta} \log [e^{\beta h} (\cosh(\beta h) + \sqrt{\sinh^2(\beta h) + e^{-4\beta^2}})]$

$$m(h, \beta) = \frac{\partial \log z}{\partial \beta h} = \frac{\sinh(\beta h) + \frac{1}{2}(\sinh^2(\beta h) + e^{-4\beta^2})^{-1/2} \cdot 2 \sinh \beta h \cosh \beta h}{\cosh(\beta h) + \sqrt{\sinh^2(\beta h) + e^{-4\beta^2}}}$$

as $h \rightarrow 0$, $\sinh(\beta h) \rightarrow 0$, $\cosh(\beta h) \rightarrow 1$

$$\text{so } \lim_{h \rightarrow 0} (m(h, \beta)) = 0$$

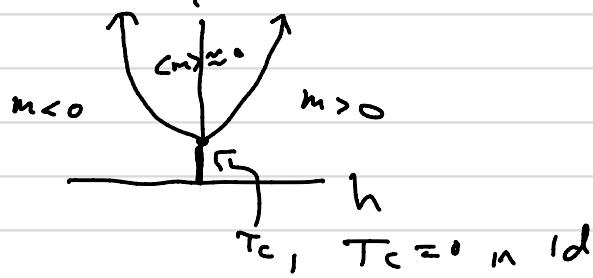
@ $\beta < \infty$,

if $\beta \rightarrow \infty$ ($T \rightarrow 0$), $e^{-4\beta^2} \rightarrow 0$ and

$$m \rightarrow \frac{\sinh(\beta h) \pm \cosh(\beta h)}{\cosh(\beta h) \pm \sinh(\beta h)} \rightarrow \pm 1$$

and which depends on $h > 0$ or $h < 0$

so phase transition / critical point at $T \rightarrow 0$, then $h \rightarrow 0$



What about higher dimensions? We can solve approx

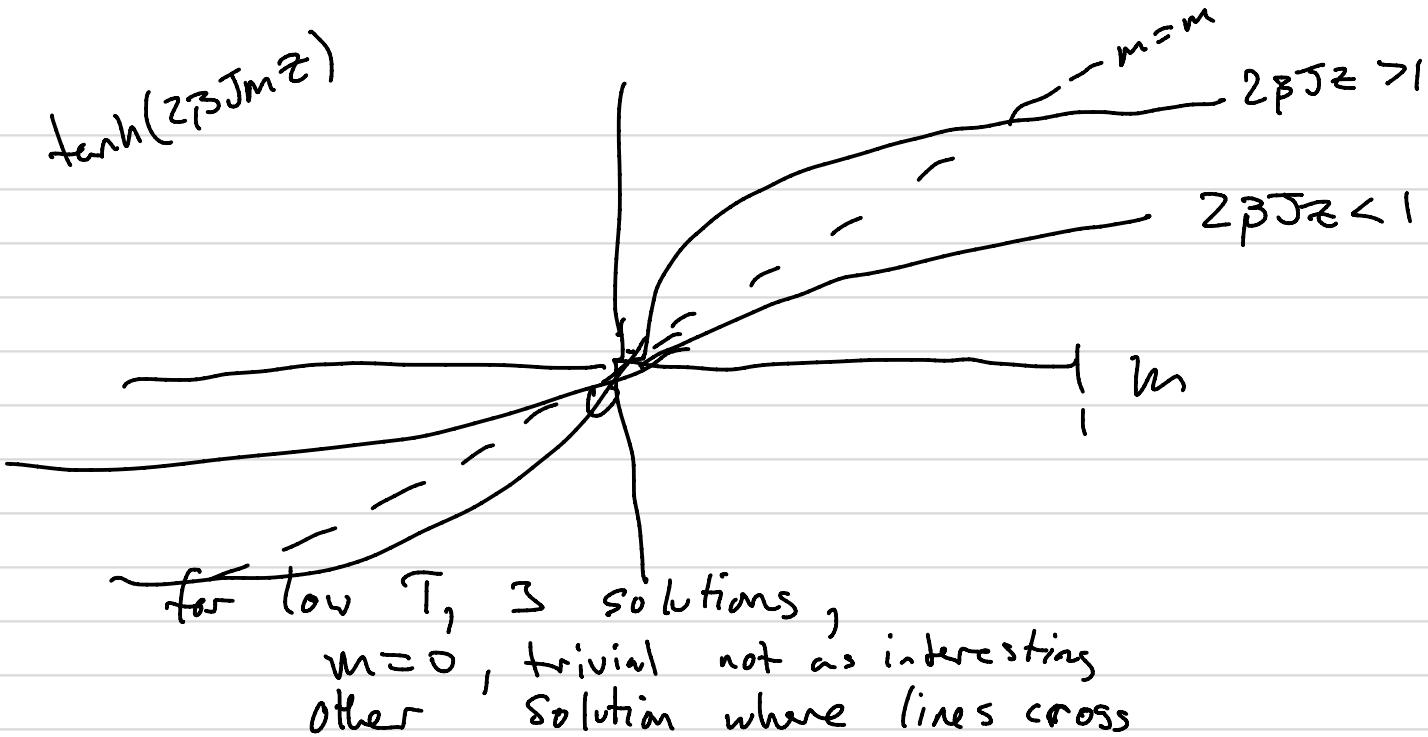
using a technique called "mean field theory", meaning each spin feels the avg effect of its neighbors

$$\text{lets call } \delta s_i = s_i - m \Rightarrow s_i = s_i + m$$

$$\begin{aligned}
 \text{then } H &= -\frac{\beta}{2} \sum_{\langle i,j \rangle} s_i s_j - h \sum_i s_i \\
 &= -\frac{\beta}{2} \sum_{\langle i,j \rangle} (m + \delta s_i)(m + \delta s_j) - h \sum_i s_i \\
 &\quad m^2 + m(\delta s_i + \delta s_j) + \overbrace{\delta s_i \delta s_j}^{\text{neglect}} \\
 &= -\frac{\beta}{2} \sum_{\langle i,j \rangle} (m^2 + m(s_i - m) + m(s_j - m)) - h \sum_i s_i \\
 &= +\frac{\beta}{2} \sum_{\langle i,j \rangle} (m^2) - \frac{-m\beta}{2} \sum_{\langle i,j \rangle} (s_i + s_j) - h \sum_i s_i \\
 &= \beta m^2 N z - (h + 2m\beta z) \sum_{i=1}^N s_i \\
 \text{So } Z &= \sum_{s_1, s_2, \dots, s_N} e^{-\beta \sum_{i=1}^N s_i} \cdot C \\
 &= e^{-\beta \sum_{i=1}^N s_i} [2 \cosh((h + 2m\beta z)\beta)]^N \\
 \langle m \rangle &= \frac{1}{N} \frac{\partial \log Z}{\partial h} = \frac{\partial \ln(2 \cosh((h + 2m\beta z)\beta))}{\partial h} = \frac{\sinh((h + 2m\beta z)\beta)}{\cosh((h + 2m\beta z)\beta)} \\
 m &= \tanh((h + 2m\beta z)\beta)
 \end{aligned}$$

No analytical solution, can get numerically
 @ $h=0$, spontaneous mag?

$$\tanh(2\beta Jmz)$$



$2\beta Jz = 1$ separates the regimes, & hence

$k_B T_c = 2\beta Jz$ gives predicted
critical point

For 1d, wrong, no $T_c > 0$