

Other Ensembles:

Constant Pressure &

Constant Chemical Potential

So far we mainly discussed situations @ const $N, V, & T$ (which is very important). However many expts actually done @

const P :



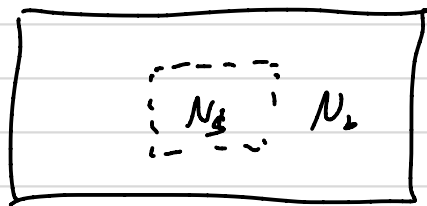
important \rightarrow density/properties depend on external pressure

Additionally, N isn't always constant, especially when there are chemical reactions

[changing # of species A, B, \dots and also maybe then absorb from gas, eg]

How do we deal with this?

Similarly



$$N = N_s + N_b$$

equil @ equal μ

$$S_b(N_b, v_b, E_b) \approx S_b(N) + \left(\frac{\partial S_b}{\partial N_b} \right) (N_b - N) + \dots$$

$-\mu/T \quad -N_s/T$

$$\Omega_b \propto e^{+\beta \mu N} [\mu, v, E]$$

Now combine transformation from $E \rightarrow T$ w/ these

to get $\Omega = P(N, P, T) \propto e^{-\beta E - \beta P V}$

$$\Delta(N, P, T) = \frac{1}{v_0} \int_0^\infty dv \int_{-\infty}^\infty dp \int_0^\infty dE e^{-\beta(\mathcal{H}(P, E) + PV)}$$

$$\approx \frac{1}{v_0} \int_0^\infty dv e^{-\beta P v} \Omega(N, v, T)$$

$$G = -k_B T \log \Delta(N, P, T) \quad \begin{array}{l} \text{Isothermal} \\ \text{Isobaric} \end{array}$$

Grand canonical:

$$\Omega_b \rightarrow P(\mu, v, T) \propto e^{-\beta E + \beta \mu N} \quad [\text{note, } \mu \text{ can be pos or neg}]$$

$$\Sigma = \sum_{N=0}^{\infty} e^{+\beta \mu N} \Omega(N, v, T)$$

$$\Omega(\mu, v, T) = -k_B T \log \Sigma(\mu, v, T) = -PV$$

(grand potential)

[Note; $e^{\beta \mu}$ sometimes called λ or activity, $\Delta \mu = k_B T \ln \lambda$]

These other ensembles are interesting because they allow for fluctuations in the quantities (U for isobaric, N for Grand) meaning the system is "compressible"

$$\frac{\epsilon_s}{\text{}} \langle N \rangle = k_B T \left(\frac{\partial}{\partial \mu} \log \mathcal{Z}(\mu, \nu, T) \right)$$

$$\text{and } \frac{\partial}{\partial \mu} = \frac{\partial \lambda}{\partial \mu} \frac{\partial}{\partial \lambda} = \beta \lambda \frac{\partial}{\partial \lambda}$$

$$\text{so } \langle N \rangle = \lambda \frac{\partial}{\partial \lambda} \log \mathcal{Z}(\lambda, \nu, T)$$

$$\sqrt{\frac{h^2}{2\pi m}} \cdot P^{1/2}$$

Ideal Gas recall $Q(N, \nu, T) = \frac{1}{N!} \left(\frac{V}{\Lambda^3} \right)^N$, $\Lambda = \sqrt{\frac{h^2}{2\pi m k_B T}}$

$$\text{so } \mathcal{Z} = \sum_{N=0}^{\infty} \frac{1}{N!} \left(\frac{V}{\Lambda^3} \right)^N \lambda^N = \exp\left(\frac{V\lambda}{\Lambda^3}\right)$$

$$PV = k_B T \log \mathcal{Z} = k_B T \cdot \frac{V\lambda}{\Lambda^3} = k_B T \langle N \rangle$$

$$\text{but } \langle N \rangle = \lambda \frac{\partial}{\partial \lambda} \left(\log \left(\exp\left(\frac{V\lambda}{\Lambda^3}\right) \right) \right) = \frac{V\lambda}{\Lambda^3}$$

$$\left[\frac{\partial \lambda}{\partial P} = \frac{1}{2} P^{-1/2} \sqrt{\frac{h^2}{2\pi m}} \right]$$

$$= \frac{1}{2} P^{-1} \Lambda$$

$$\text{so } \boxed{PV = \langle N \rangle RT}$$

$$\text{and } \langle E \rangle = -\frac{\partial}{\partial \beta} \log \mathcal{Z} = -\frac{\partial}{\partial \beta} \left(\frac{V\lambda}{\Lambda^3} \right) = V\lambda \cdot 3\lambda^{-4} \frac{\partial \lambda}{\partial \beta} = \frac{3}{2} \langle N \rangle k_B T$$

(sackr technique - exercise?)

What about fluctuations now

$$\sigma_N^2 = \langle N^2 \rangle - \langle N \rangle^2$$

Lets start w/ $\langle N^2 \rangle$

$$= \left\langle \sum_N N^2 P(N) \right\rangle_E = \left\langle \sum_N N^2 \frac{e^{-\beta E + \mu N \beta}}{Z} \right\rangle_E$$

$$= \frac{1}{Z} \left\langle \sum_N N \frac{1}{\beta} \frac{\partial}{\partial \mu} e^{\mu N \beta} \right\rangle = \frac{k_B T}{Z} \left\langle \frac{\partial}{\partial \mu} \sum_N N e^{\mu N \beta} \right\rangle_E$$

$$\langle N \rangle = \left\langle \sum_N e^{\mu N \beta} \cdot N / Z \right\rangle_E \Rightarrow \langle N \rangle Z = \left. \right$$

$$\begin{aligned} \langle N^2 \rangle &= \frac{k_B T}{Z} \frac{\partial}{\partial \mu} (\langle N \rangle Z) = \frac{k_B T}{Z} \cdot \left(\langle N \rangle \frac{\partial}{\partial \mu} Z + Z \frac{\partial}{\partial \mu} \langle N \rangle \right) \\ &= \langle N^2 \rangle + k_B T \frac{\partial}{\partial \mu} \langle N \rangle \end{aligned}$$

$$\Rightarrow k_B T \frac{\partial}{\partial \mu} \langle N \rangle = \langle N^2 \rangle - \langle N \rangle^2 = \sigma_N^2 \quad \text{FDT}$$

$$\langle N \rangle = k_B T \frac{\partial}{\partial \mu} \log Z = \frac{\partial}{\partial \mu} (k_B T \log Z)$$

$$\text{so } \sigma_N^2 = (k_B T)^2 \frac{\partial^2}{\partial \mu^2} \log Z = k_B T V \frac{\partial^2 P}{\partial \mu^2}$$

non-trivial

Non-trivial to make this useful

$$a(T, \nu) = \frac{F(N, V, T)}{N}, \quad \nu = V/N \quad (\text{only depends on intensive variables})$$

$$\mu = \frac{\partial F}{\partial N} = \frac{\partial}{\partial N} (aN) = a + \frac{\partial a}{\partial \nu} \nu = a + N \left(\frac{\partial a}{\partial (V/N)} \right) \frac{\partial (V/N)}{\partial N} = a - \nu \frac{\partial a}{\partial \nu}$$

$$[dF = -SdT - PdV + \mu dN]$$

$$\frac{\partial \mu}{\partial v} = \frac{\partial}{\partial v} \left(a - v \frac{\partial a}{\partial v} \right) = \frac{\partial a}{\partial v} - v \frac{\partial^2 a}{\partial v^2} - \frac{\partial a}{\partial v} = -v \frac{\partial^2 a}{\partial v^2}$$

Pressure

$$P = - \frac{\partial F}{\partial V} = - \frac{\partial (NF/v)}{\partial v} = -N \frac{\partial a}{\partial v} \frac{\partial (v/v)}{\partial v} = - \frac{\partial a}{\partial v}$$

$$\frac{\partial P}{\partial v} = - \frac{\partial^2 a}{\partial v^2} \Rightarrow \boxed{\frac{\partial \mu}{\partial v} = v \frac{\partial P}{\partial v}}$$

Also: $\left(\frac{\partial P}{\partial \mu} \right)_{T,N} = \frac{\partial}{\partial \mu} \left(\frac{k_B T}{v} \ln z \right) = \frac{1}{v} \frac{\partial}{\partial \mu} (k_B T \ln z) = \langle N \rangle / v = \frac{1}{v}$

Back to $\sigma_N^2 = k_B T v \frac{\partial^2 P}{\partial \mu^2} = k_B T v \frac{\partial}{\partial \mu} \left(\frac{\partial P}{\partial \mu} \right)$

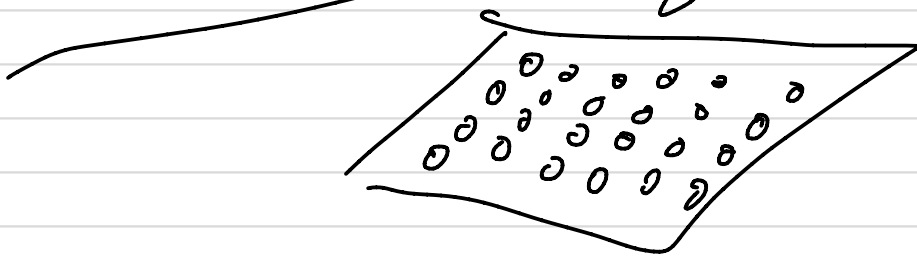
$$= k_B T v \left(\frac{\partial v}{\partial \mu} \right) \frac{\partial}{\partial v} \left(\frac{\partial P}{\partial v} \right) = k_B T v \cdot \frac{1}{v^2} \cdot \frac{\partial v}{\partial \mu}$$

$$= k_B T v \cdot \frac{1}{v^2} \cdot \left(\frac{1}{v} \frac{\partial v}{\partial P} \right)$$

$$\boxed{\sigma_N^2 = \frac{k_B T \langle N \rangle^2}{v} \cdot K} \quad \begin{matrix} ||| \\ -K_T \end{matrix}$$

Example problem

ideal gas



@ ϵ_f when $\mu_{\text{gas}} = \mu_{\text{substrate}}$

if site occupied or unoccupied & each site has partition function q when vac & $qe^{\mu\beta}$ otherwise

$$Q(N, M, T) = \binom{M}{N} q^N$$

$$Z = \sum_{N=0}^M e^{\mu N \beta} Q(N) = \sum_{N=0}^M \binom{M}{N} (e^{\mu\beta} q)^N (1)^{M-N}$$
$$= (1 + qe^{\mu\beta})^M$$

$$\langle N \rangle = k_B T \frac{\partial \log Z}{\partial \mu} = k_B T M \frac{\partial}{\partial \mu} \log(1 + qe^{\mu\beta})$$

$$= k_B T M \cdot \frac{qe^{\mu\beta}}{1 + qe^{\mu\beta}}$$

$$\langle N \rangle / M = \theta = \frac{qe^{\mu\beta}}{1 + qe^{\mu\beta}}$$

For the gas $\mu = \left(\frac{\partial A}{\partial N} \right)_{T, V}$, $A = -k_B T \ln Z$

$$\mu_g = -k_B T \frac{\partial}{\partial N} \left(\ln \left(\frac{1}{N!} \left(\frac{V}{\Lambda^3} \right)^N \right) \right)$$

$z_g \downarrow$
 $z = \frac{1}{N!} \left(\frac{V}{\Lambda^3} \right)^N$
 $\uparrow q_g$

$$= -k_B T \frac{\partial}{\partial N} (N \ln q_g - N \ln N + N)$$

$$= -k_B T (\ln q_g + 1 - (1 + \ln N))$$

$$PV = N k_B T$$

$$= -k_B T \ln \left(\frac{q_g}{N} \right)$$

$$\Theta = \frac{q_s / \left(\frac{q_g}{N} \right)}{1 + \left(\frac{q_s}{\left(\frac{q_g}{N} \right)} \right)} = \frac{\left(\frac{q_g}{N} \right)^{-1}}{\frac{1}{q_s} + \left(\frac{q_g}{N} \right)^{-1}}$$

$$q_g = \frac{V}{\Lambda^3} \Rightarrow \frac{q_g}{N} = \left(\frac{V}{N} \right) \cdot \frac{1}{\Lambda^3} = \frac{k_B T}{P \Lambda^3}$$

$$= \frac{P \Lambda^3 \beta}{\frac{1}{q_s} + P \Lambda^3 \beta} = \frac{P}{P + \frac{1}{\Lambda^3 \beta q_s}}$$

$$= \frac{P}{P + P_0}$$

$$P_0 = \frac{k_B T}{q_s} \left(\frac{2 \pi m k_B T}{h^2} \right)^{3/2}$$

