

HW2 - Glen Hockey

① Gamma functions

$$\Gamma(z) = \int_0^\infty x^{z-1} e^{-x} dx$$

a) $\Gamma(1) = \int_0^\infty x^0 e^{-x} dx = \int_0^\infty e^{-x} dx$

$$= -e^{-x} \Big|_0^\infty = 0 - (-1) = 1 \quad \checkmark$$

b) $\Gamma(z+1) = \int_0^\infty x^z e^{-x} dx$

$$= \underbrace{-x^z e^{-x}}_{0-0} \Big|_0^\infty + \int_0^\infty e^{-x} \cdot z x^{z-1} dx = z \Gamma(z) \quad \checkmark$$

$\left[\int u dv = uv - \int v du \right]$

$u = x^z \quad du = z x^{z-1}$
 $dv = e^{-x} \quad \Rightarrow v = -e^{-x}$

c) $\Gamma(N+1) = N \Gamma(N) = N \cdot (N-1) \Gamma(N-1) = \dots$

$$= N(N-1)(N-2) \dots \Gamma(1)$$

$$= N(N-1)(N-2) \dots (1) = N! \quad \checkmark$$

2 a) for $d=2$,

$$V_2 = \pi r^2, \quad S_1 = 2\pi r$$

$$\frac{S_1}{V_2} = 2 = d \quad \checkmark$$

$$V_3 = 4/3\pi r^3 \quad S_2 = 4\pi r^2$$

$$\frac{S_2}{V_3} = 3 = d \quad \checkmark$$

$$b)i) I = \int_{-\infty}^{\infty} dx e^{-x^2}$$

$$\begin{aligned} I^d &= \pi \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} dx_1 dx_2 e^{-x_1^2 - x_2^2} \\ &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \cdots dx_1 dx_2 \cdots e^{-\sum_{i=1}^d x_i^2} \end{aligned}$$

$$I^d = \int_{-\infty}^{\infty} dx^d e^{-r^2} = S_{d-1} \int_0^{\infty} r^{d-1} e^{-r^2} dr$$

$$\text{ii) } \int_0^{\infty} r^{d-1} e^{-r^2} dr = \int_0^{\infty} dy r^{d-2} e^{-y} = \int_0^{\infty} dy y^{\frac{d-2}{2}} e^{-y}$$

$\uparrow y = r^2 \quad dy = 2r dr \Rightarrow dr = \frac{dy}{2r}$

$$= \frac{1}{2} \int_0^{\infty} dy e^{-y} y^{\frac{d-2}{2}-1} = \frac{1}{2} \Gamma(d/2)$$

$$\text{iii) } I = \sqrt{\pi} \quad \text{so} \quad I^d = \pi^{d/2}$$

$$\Rightarrow \pi^{d/2} = \frac{1}{2} \Gamma(d/2) \cdot S_{d-1}$$

$$\Rightarrow S_{d-1} = \frac{2\pi^{d/2}}{\Gamma(d/2)}$$

$$iv) \quad V_d = S_{d-1}/d$$

$$\Rightarrow V_d = \frac{2\pi^{d/2}}{d \Gamma(d/2)} = \frac{\pi^{d/2}}{(d/2)\Gamma(d/2)} = \frac{\pi^{d/2}}{\overbrace{\Gamma(d/2+1)}} \quad \sim$$

$$c) i) \quad \Gamma(t) = \int_0^\infty dx e^{-x} x^{t-1}$$

$$\Gamma(k) = \int_0^\infty dx e^{-x} x^{-k} = \int_0^\infty r dr e^{-r^2} r^{-1} =$$

$$\stackrel{?}{=} x = r^2, \quad dx = 2r dr$$

$$= 2 \int_0^\infty dr e^{-r^2} = \int_{-\infty}^\infty dr e^{-r^2} = \sqrt{\pi} \quad \approx$$

$$ii) \quad V_3 = \pi^{3/2} / \Gamma(3/2 + 1)$$

$$\rightarrow \Gamma(3/2 + 1) = 3/2 \Gamma(1/2 + 1) = 3/2 \cdot 1/2 \Gamma(1/2) = \frac{3}{4} \sqrt{\pi}$$

$$\Rightarrow V_3 = 4/3 \pi \quad \checkmark$$

remember
 $\Gamma(N+1) \approx N!$

3) Microcanonical ideal gas

$$S(N, V, \epsilon) = \frac{\epsilon_0 V^N}{h^{3N} N!} \int_{-\infty}^{\infty} dp^{3N} S \left(\sum_{i=1}^{3N} p_i^2 - \epsilon \right)$$

$p_i = \sqrt{2m} y_i$

$$= \frac{\epsilon_0 V^N (2m)^{3N/2}}{h^{3N} N!} \int_{-\infty}^{\infty} dy^{3N} S \left(\sum_{i=1}^{3N} y_i^2 - \epsilon \right)$$

$$\text{let } r^2 = \sum_{i=1}^{3N} y_i^2$$

a) eq 3:

$$\int dy^{3N} f(y) = S_{3N-1} \int_0^\infty dr f(r) r^{3N-1}$$

eq 4

$$= 2\pi \frac{(3N-1)/2}{\Gamma(\frac{3N}{2})} \int_0^\infty dr f(r) r^{3N-1}$$

so

$$S(N, V, \epsilon) = \frac{\epsilon_0 V^N (2m)^{3N/2}}{h^{3N} N!} \cdot \frac{2\pi^{(3N/2)}}{\Gamma(\frac{3N}{2})} \int_0^\infty dr S(r^2 - \epsilon) r^{3N-1}$$

b)

$$\int_0^\infty \delta(r^2 - \varepsilon) r^{3N-1} = \frac{1}{2\sqrt{\varepsilon}} \cdot \left(\int_0^\infty \delta(r - \sqrt{\varepsilon}) r^{3N-1} + \int_0^\infty \delta(r + \sqrt{\varepsilon}) r^{3N-1} \right)$$

$\approx 0, -\sqrt{\varepsilon}$ outside range

so

$$\int_0^\infty dr \delta(r^2 - \varepsilon) r^{3N-1} = \frac{1}{2\sqrt{\varepsilon}} \int_0^\infty \delta(r - \sqrt{\varepsilon}) r^{3N-1} dr$$

$$\approx \frac{1}{2\sqrt{\varepsilon}} (\sqrt{\varepsilon})^{3N-1} = \frac{1}{2\varepsilon} \cdot \varepsilon^{(3N)/2}$$

therefore

$$\begin{aligned} S(N, V, \varepsilon) &= \frac{\varepsilon_0 V^N (2m)^{3N/2}}{h^{3N} N!} \cdot \frac{2\pi^{(3N/2)}}{\Gamma(3N/2)} \cdot \frac{1}{2\varepsilon} \cdot \varepsilon^{(3N/2)} \\ &= \left(\frac{\varepsilon_0}{\varepsilon}\right) \frac{V^N (2m)^{3N/2}}{h^{3N} N!} \frac{\pi^{(3N/2)}}{\Gamma(3N/2)} \varepsilon^{3N/2} \end{aligned}$$

$$S(N, V, \varepsilon) = \frac{\varepsilon_0}{\varepsilon} \frac{1}{N! \Gamma(3N/2)} \left[V \left(\frac{2\pi m \varepsilon}{h^2} \right)^{3/2} \right]^N$$

c)

$$S(N, V, \epsilon) = \frac{\epsilon_0}{\epsilon} \frac{1}{N!} \Gamma\left(\frac{3N}{2}\right) \left[V \left(\frac{2\pi m \epsilon}{h^2} \right)^{3/2} \right]^N$$

$$\Gamma\left(\frac{3N}{2}\right) = \left(\frac{3N}{2} - 1\right)! \approx \left(\frac{3N}{2}\right)! \approx \left(\frac{3N}{2}\right)^{\frac{3N}{2}} e^{-3N/2}$$

And $\epsilon^{3N/2 - 1} \approx \epsilon^{3N/2}$

$$so S(N, V, \epsilon) \approx \frac{\epsilon_0}{N!} \cdot \frac{1}{\left(\frac{3N}{2}\right)^{\frac{3N}{2}} e^{-3N/2}} \left[V \left(\frac{2\pi m \epsilon}{h^2} \right)^{3/2} \right]^N$$

$$= \frac{\epsilon_0}{N!} \cdot \left[V \cdot \left(\frac{4\pi m \epsilon e}{3h^2 N} \right)^{3/2} \right]^N$$

d) $\epsilon = \frac{3}{2} N k_B T$

$$1) so S(N, V, \epsilon) \approx \frac{\epsilon_0}{N!} \cdot \left[V \left(\frac{2\pi m k_B T}{h^2} \right)^{3/2} \right]^N \cdot e^{\frac{3N}{2}}$$

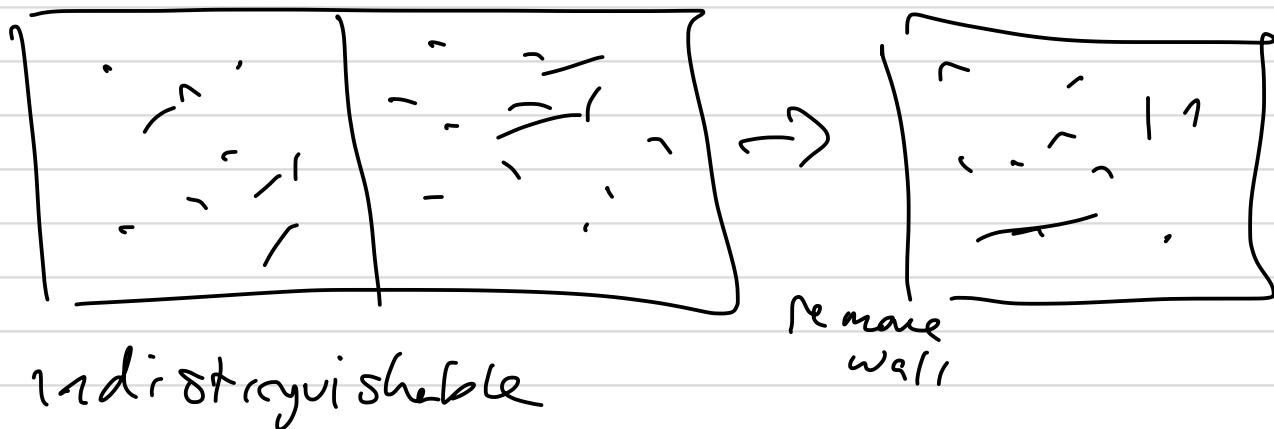
$$S = k_B \ln \mathcal{J} \approx N k_B \ln \left[V \left(\frac{2 \pi m k_B T}{n^2} \right) \right] + \frac{3N}{2} k_B - k_B \ln N! \checkmark$$

$$\text{i)} -k_B \ln N! \approx -k_B [N \ln N - N]$$

$$S \approx N k_B \ln \left[\frac{V}{N} \left(\frac{2 \pi m k_B T}{n^2} \right) \right] + \frac{5N}{2} k_B \checkmark$$

4) Gibbs paradox, entropy of mixing

What if we didn't have $1/N!$



$$S_{\text{cl}}^1 \sim N_1 k \log v_1 + 3/2 N_1 k + \text{const} \cdot N_1$$

$$S_{\text{cl}}^2 \sim N_2 k \log v_2 + 3/2 N_2 k + \text{const} \cdot N_2$$

$$S_{\text{cl}}^f = (N_1 + N_2)k \log(v_1 + v_2) + \frac{5}{2}(N_1 + N_2)k + \text{const}(N_1 + N_2)$$

$$\Delta S_{\text{cl}} = N_1 k \log \left(\frac{v}{v_1} \right) + N_2 k \log \left(\frac{v}{v_2} \right) > 0 \quad \text{since}$$

$$v > v_1 \quad \& \quad v > v_2$$

However, w/ correction

$$\begin{aligned} \Delta S &\geq (N_1 + N_2)k \log \left(\frac{v_1 + v_2}{N_1 + N_2} \right) + \frac{5}{2}(N_1 + N_2)k \\ &\quad - \left(N_1 k \log \frac{v}{v_1} + N_2 k \log \frac{v}{v_2} + \frac{5}{2}(N_1 + N_2)k \right) \end{aligned}$$

$$= N_1 k \log \frac{v}{N_1 v_1} + N_2 k \log \frac{v}{N_2 v_2} \approx 0$$

because we start at equilibrium, where

$$N/v = N_1/v_1 = N_2/v_2 = P/k_B T$$