

# Homework 3: Canonical ensemble and radial distribution function

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1. (Computational) This exercise will demonstrate how to compute a radial distribution function ( $g(r)$ ) for somewhat realistic simulation data. It is available by logging in to <https://chemga-2600-fall.rcnyu.org/> and going to the assignments tab. There you can push "Fetch" next to RDF. From here you can access the RadialDistributionFunction exercise. When done, click validate and then submit.

*Alternatively*, fill in the missing pieces in the notebook within the rdf folder on the course github:

<https://github.com/hockyg/chem-ga-2600/>

When finished, download with File->Download As->ipynb so you have a copy and send it in by email or slack.

2. *Ideal gas of molecules* (Adapted from Tuckerman problem 4.12). If you have  $N$  identical non-interacting molecules (each with  $n$  atoms) in a box, the total partition function factorizes,

$$Q(N, V, T) = \frac{q(n, V, T)^N}{N!} \quad (1)$$

Moreover, the single molecule partition function  $q(n, V, T) = Vf(n, T)$ , where  $f$  is a function that only depends on the number of atoms and the temperature.

Now suppose that the system contains different types of molecules, this factorization still works. For example, with 2 types,  $A$  and  $B$ ,

$$Q(N_A, N_B, V, T) = \frac{q_A(n_A, V, T)^{N_A}}{N_A!} \cdot \frac{q_B(n_B, V, T)^{N_B}}{N_B!} \quad (2)$$

- (a) An example chemical reaction might be



The Helmholtz free energy  $A$  is now a function of  $V, T$ , and all 4  $N$ 's. Let there be a variable called  $\lambda$  which is the reaction extent. Then  $dN_A = a d\lambda$ ,  $dN_B = b d\lambda$ ,  $dN_C = -c d\lambda$ , and  $dN_D = -d d\lambda$ . At chemical equilibrium,  $A$  is at a minimum, which means that  $dA/d\lambda = 0$ .

**Show that at equilibrium with fixed  $V$  and  $T$ ,**

$$a\mu_A + b\mu_B - c\mu_C - d\mu_D = 0, \quad (4)$$

where chemical potentials of each species are defined like:

$$\mu_A = -k_B T \frac{\partial \ln Q(V, T, N_A, N_B, N_C, N_D)}{\partial N_A} \quad (5)$$

Hint: This should follow very directly from writing out the chain rule for  $dA$ .

- (b) By plugging in for  $Q$  in the 4 equations like Eq. 5 and substituting this in to Eq. 4, show that you get the following relationship

$$K(T) = \frac{\rho_C^c \rho_D^d}{\rho_A^a \rho_B^b} = \frac{(q_C/V)^c (q_D/V)^d}{(q_A/V)^a (q_B/V)^b}, \quad (6)$$

where  $\rho_A = N_A/V$ , etc. Hint, you may have to use Sterling's approximation. Do you see why both the middle and right fractions are functions of only temperature?

- (c) Using the definition of the Helmholtz free energy and the formula for  $Q$  with 4 different species, show that the Helmholtz free energy  $A$  can be written as  $A_{total} = A_A + A_B + A_C + A_D$ . What is the equation for  $A_X$ , that is, the contribution to the Helmholtz free energy from species  $X$ .
- (d) The pressure in the canonical ensemble can be obtained by the formula  $P = -\left(\frac{\partial A}{\partial V}\right)_{N,T}$ . The partial pressure of species  $X$  can be expressed similarly as  $P_X = -\left(\frac{\partial A_X}{\partial V}\right)_{N,T}$ , where  $A_X$  was defined in the previous part. Given this definition, find the relationship between  $K(T)$  above and

$$K_P = \frac{P_C^c P_D^d}{P_A^a P_B^b} \quad (7)$$

3. *Deriving the Canonical Ensemble from the perspective of Maximum Entropy.* Consider a system with  $N$  possible states each having energy  $\epsilon_i$ . One definition we read for the entropy is the Gibbs entropy,  $S = -k_B \sum_{i=1}^N p_i \ln p_i$ , where  $p_i$  is the (unknown) probability that system is in state  $p_i$ .

By definition,  $\sum_{i=1}^N p_i = 1$ , and the average energy  $\bar{E}$  is a constant, with  $\bar{E} = \sum_{i=1}^N \epsilon_i p_i$ . Note that the Gibbs entropy is strictly positive since  $0 < p_i < 1 \Rightarrow -\ln p_i > 0$ .

A question we could ask is, which set of  $\{p_i\}$  maximize the Gibbs entropy subject to the constraints  $\sum_{i=1}^N p_i = 1$  and  $\bar{E} = \sum_{i=1}^N \epsilon_i p_i$ . To do this, we can use the method of undetermined Lagrange Multipliers. If you don't know what this is, you should read about it a little on your own. The upshot is, we have to maximize the quantity:

$$L = -k_B \sum_{i=1}^N p_i \ln p_i - \alpha \sum_{i=1}^N p_i - \beta \sum_{i=1}^N \epsilon_i p_i, \quad (8)$$

where  $\alpha$  and  $\beta$  are undetermined variables.

- (a) Find where  $L$  is maximum as a function of an arbitrary  $p_k$  by solving  $\frac{dL}{dp_k} = 0$  for  $p_k$ , where the  $p$ 's are considered all independent variables.
- (b) Use the fact that  $\sum_{i=1}^N p_i = 1$  to eliminate one of the unknown variables. Your expression for the probabilities should look familiar.
- (c) Based on what you know about the Canonical ensemble, express this  $\beta$  in terms of a familiar thermodynamic variable.
- (d) Optional extra part: what happens if we don't include the constraint on the average energy? What ensemble does this correspond to?